

# Diffraction and Exclusive Production at CDF II



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**LISHEP 2009**

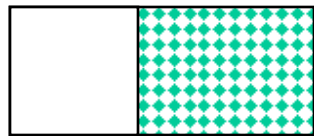
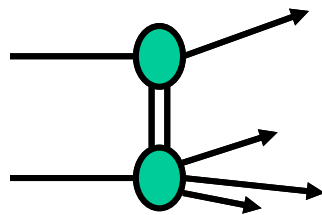
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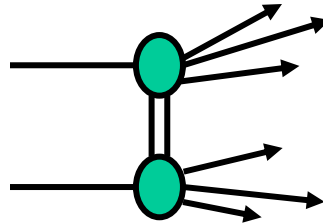
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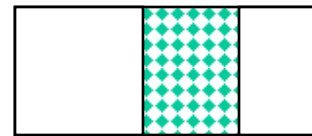
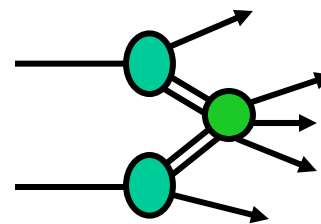
# Soft and hard diffractive and exclusive studies at CDF



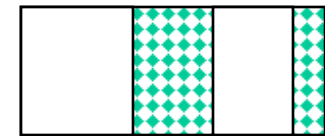
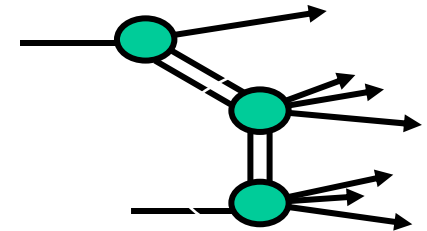
SD



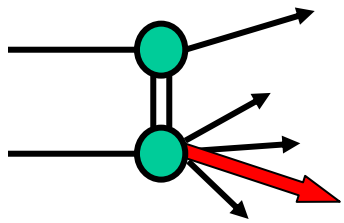
DD



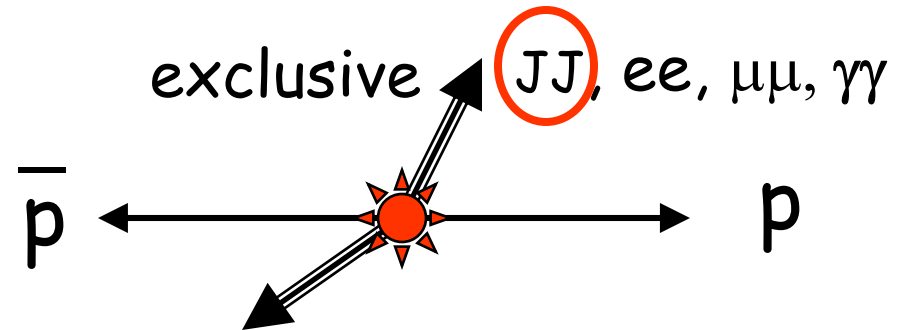
DPE

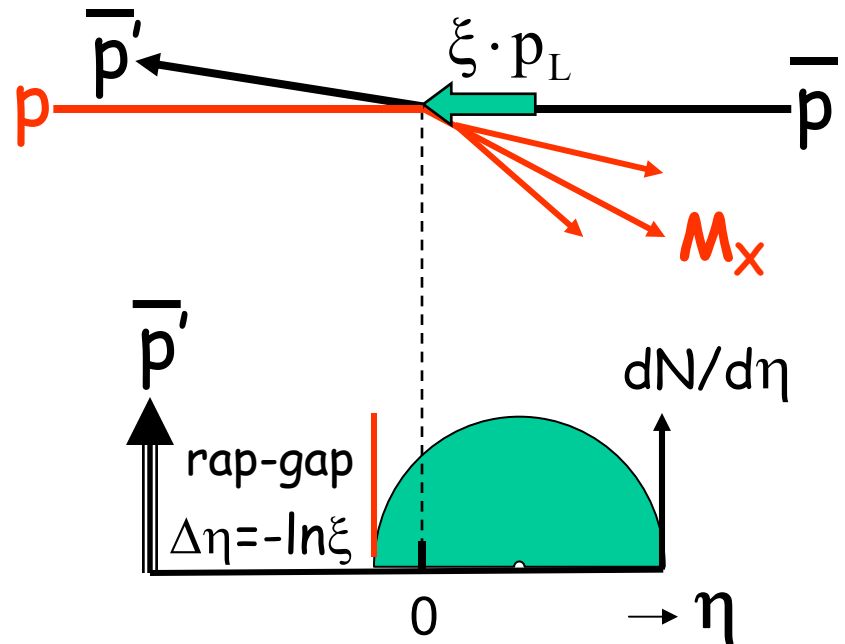
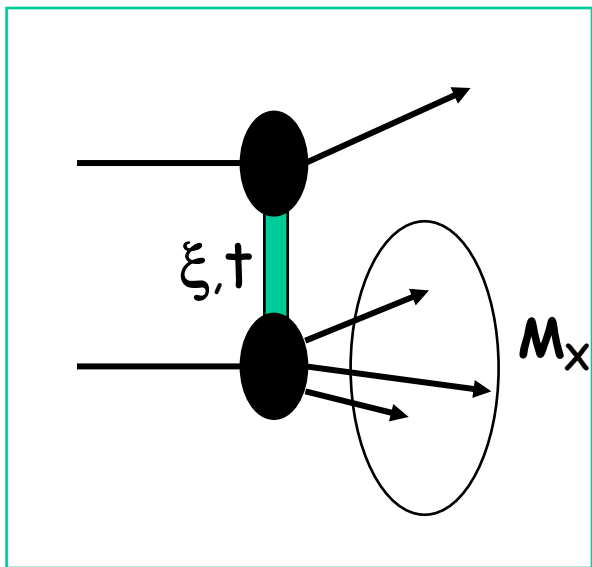
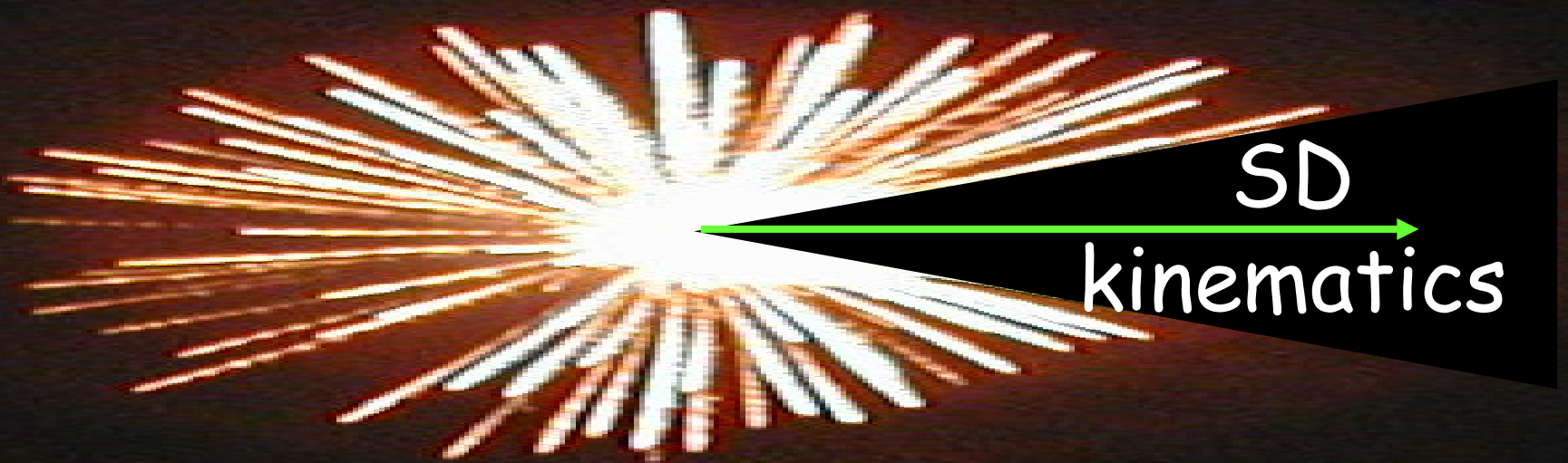


SDD=SD+DD

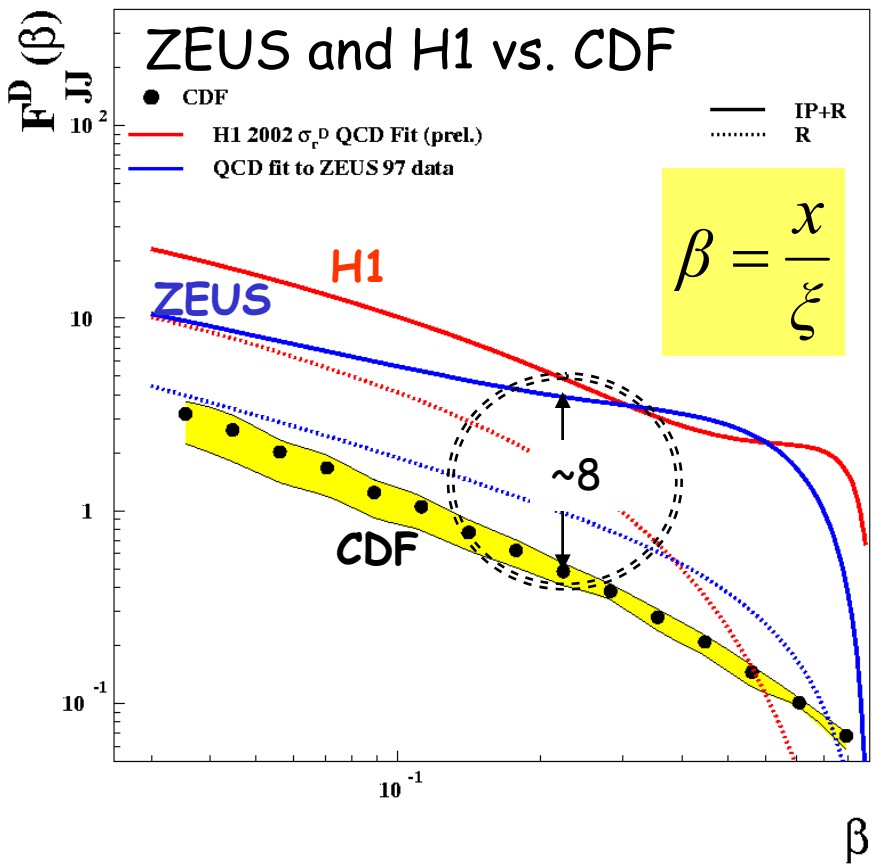
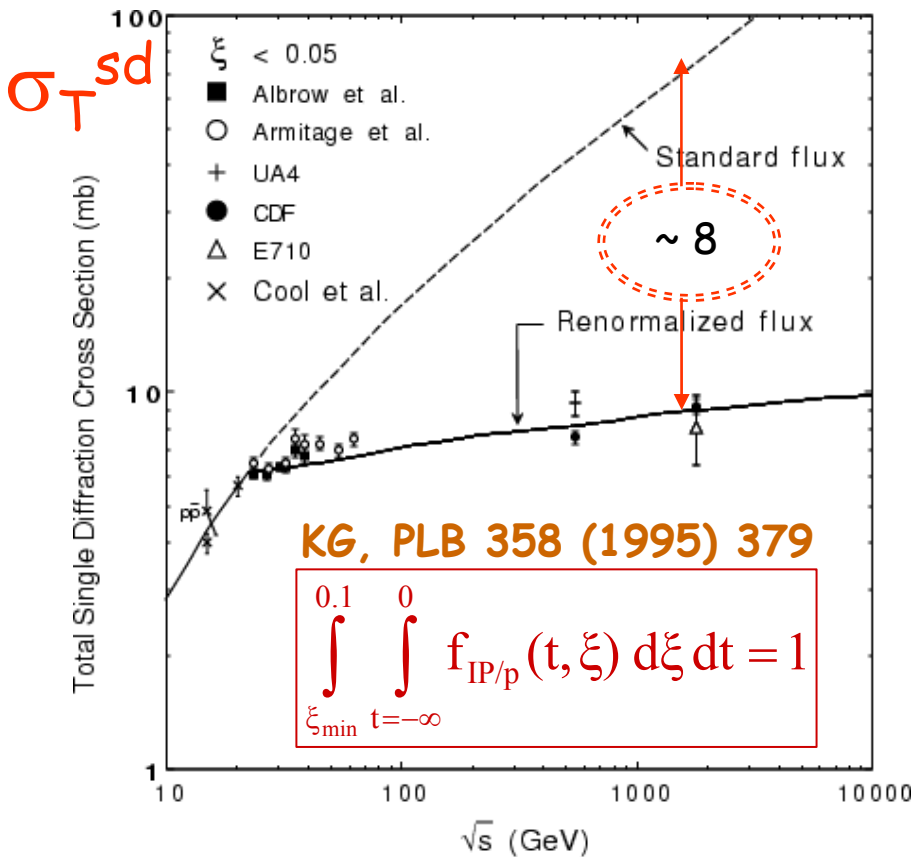
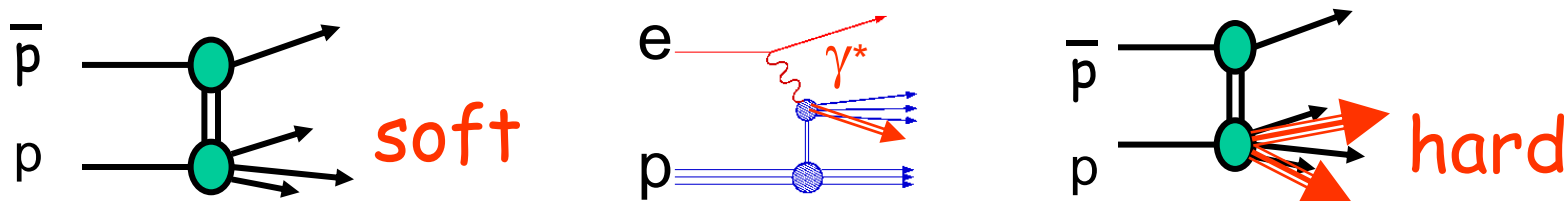


JJ, b, J/ψ, W





# Breakdown of factorization



**Magnitude:** same suppression factor in soft and hard diffraction!  
**Shape of  $\beta$  distribution:** ZEUS, H1, and Tevatron - why different slapes?

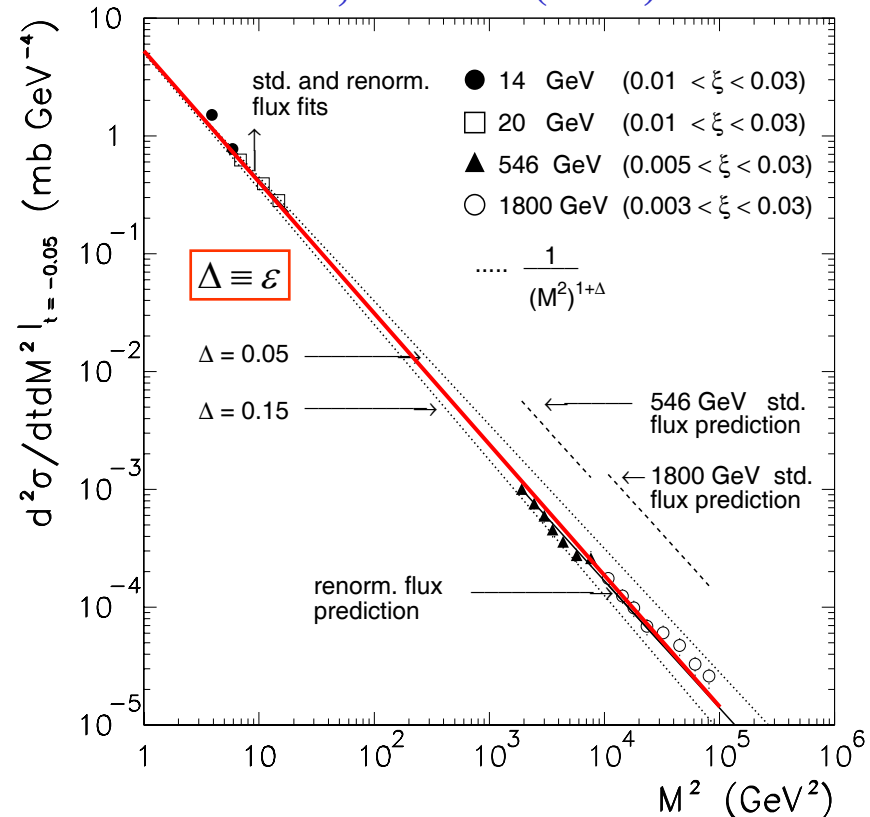
→ ds/dM<sup>2</sup> independent of s !

renormalization

$$\frac{d\sigma}{dM^2} \propto \frac{s^{2\varepsilon} \rightarrow 1}{(M^2)^{1+\varepsilon}}$$

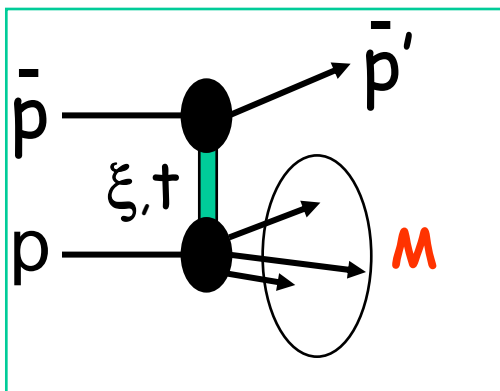
→ Independent of s over 6 orders of magnitude in M<sup>2</sup>!

KG&JM, PRD 59 (1999) 114017

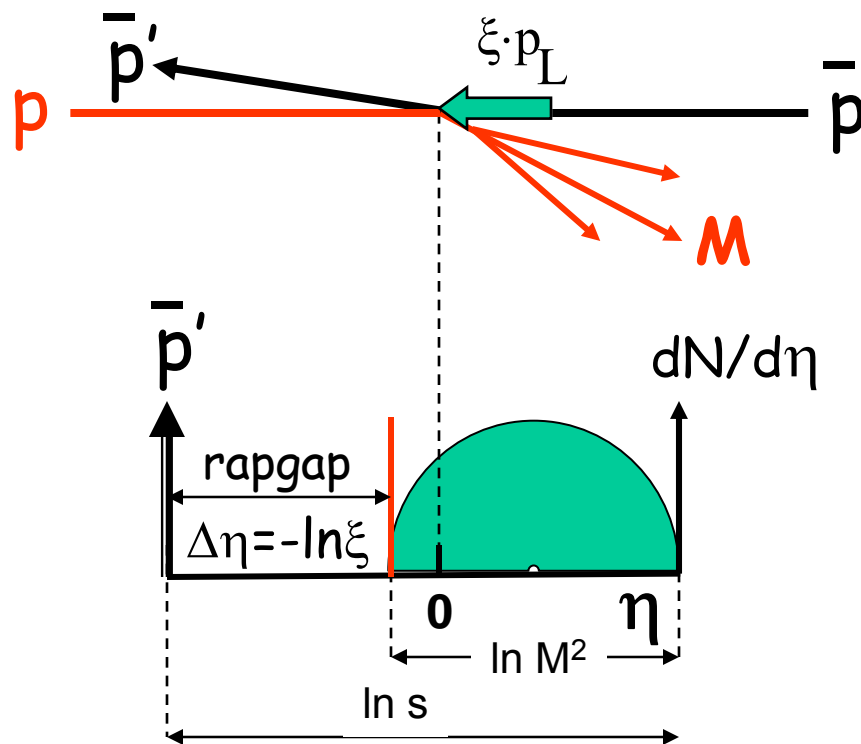


Factorization breaks down so as to ensure M<sup>2</sup> scaling!

$M^2$  scaling  
expected in QCD



$$1 - x_L \equiv \xi = \frac{M^2}{s}$$



**vacuum exchange**




$$\left( \frac{d\sigma}{d\Delta\eta} \right)_{t=0} \approx \text{constant} \Rightarrow \frac{d\sigma}{d\xi} \propto \frac{1}{\xi} \Rightarrow \frac{d\sigma}{dM^2} \propto \frac{1}{M^2}$$

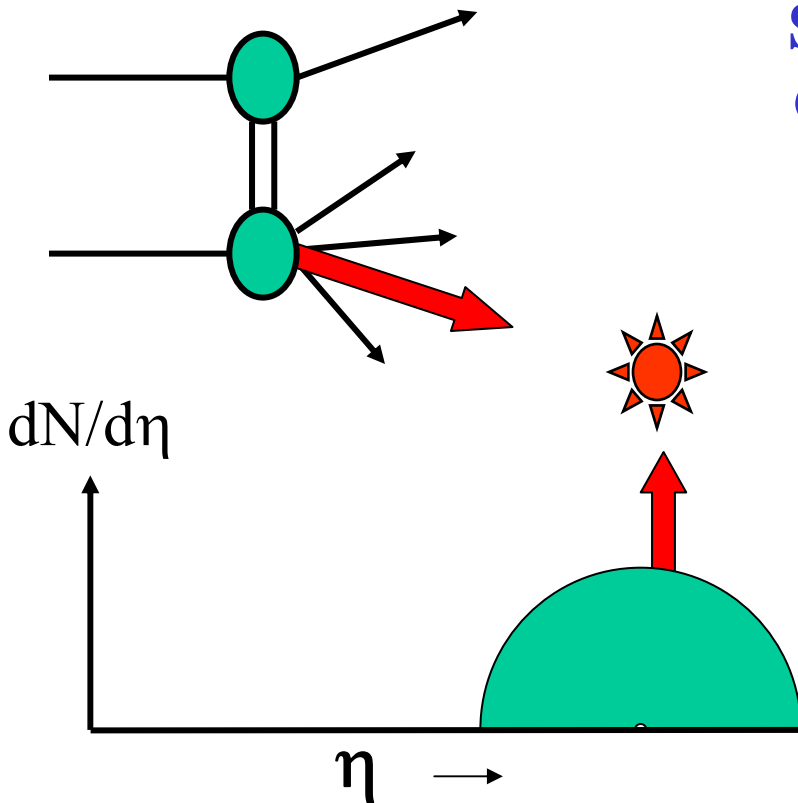


# Hard diffractive fractions

$$\bar{p}p \rightarrow (\text{Sun} + X) + \text{gap}$$

Fraction:  
SD/ND ratio  
@ 1800 GeV

	Fraction %
JJ	0.75 +/- 0.10
W	0.115 +/- 0.55
b	0.62 +/- 0.25
J/ψ	1.45 +/- 0.25

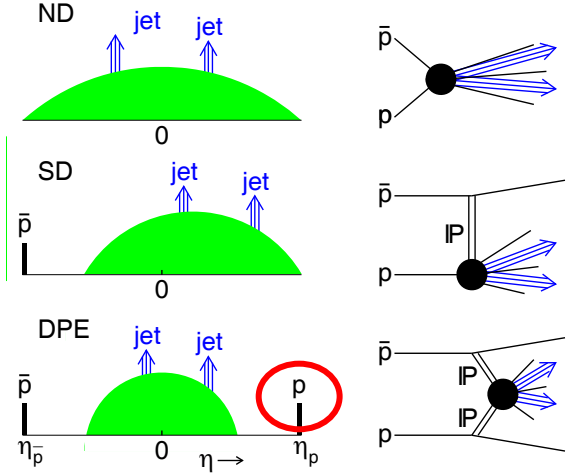


All fractions ~ 1%  
(differences due to kinematics)

- ~ uniform suppression
- ~ **FACTORIZATION ! !**



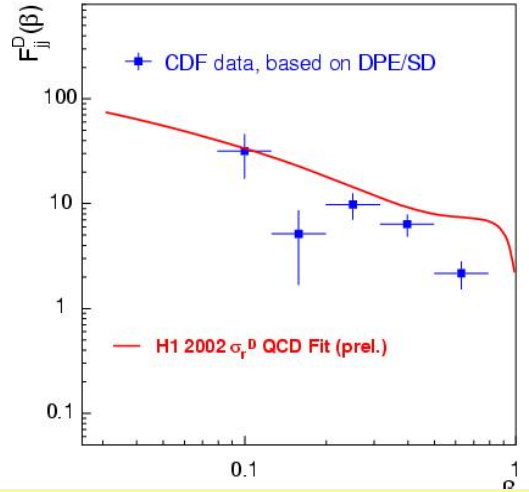
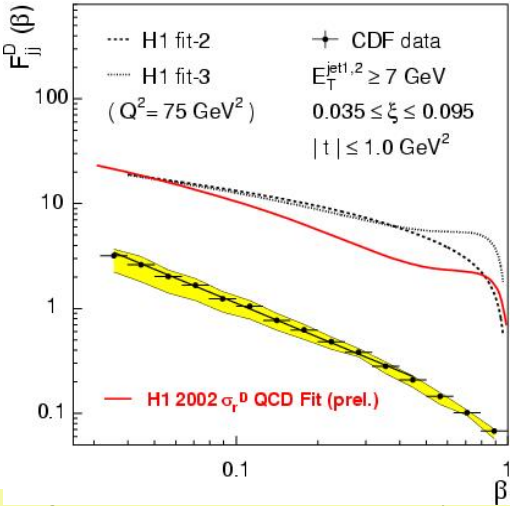
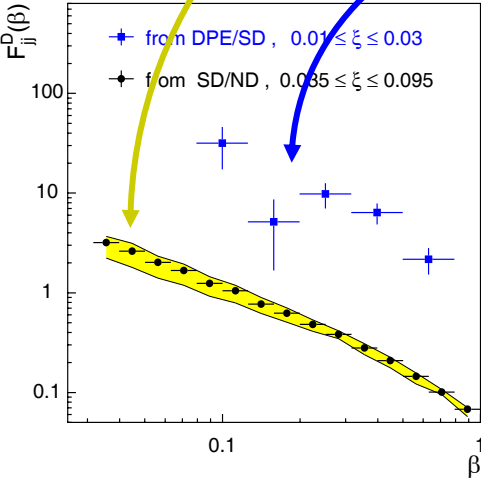
# Multi-gap diffraction - restoring factorization -



$R(\text{SD}/\text{ND})$

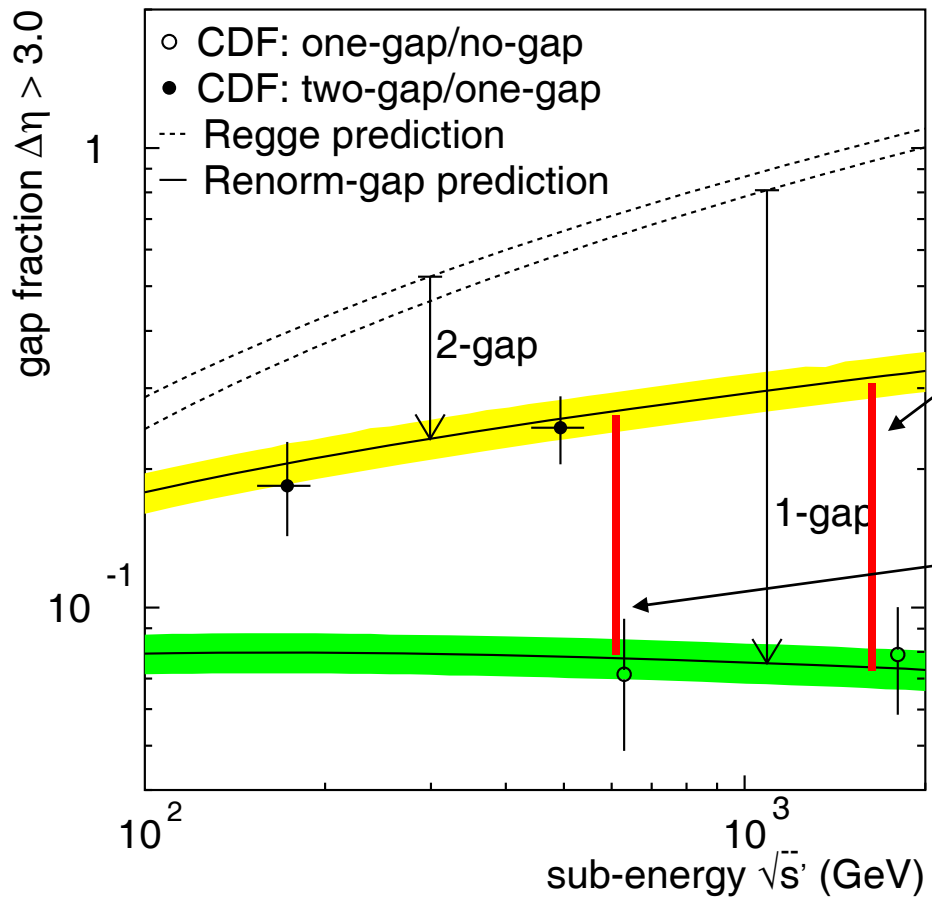
$R(\text{DPE}/\text{SD})$

DSF from ratio of two/one gap:  
factorization restored!



The diffractive structure function fit measured on the proton side in events with a leading antiproton is NOT suppressed relative to predictions based on DDIS

# Gap Survival Probability

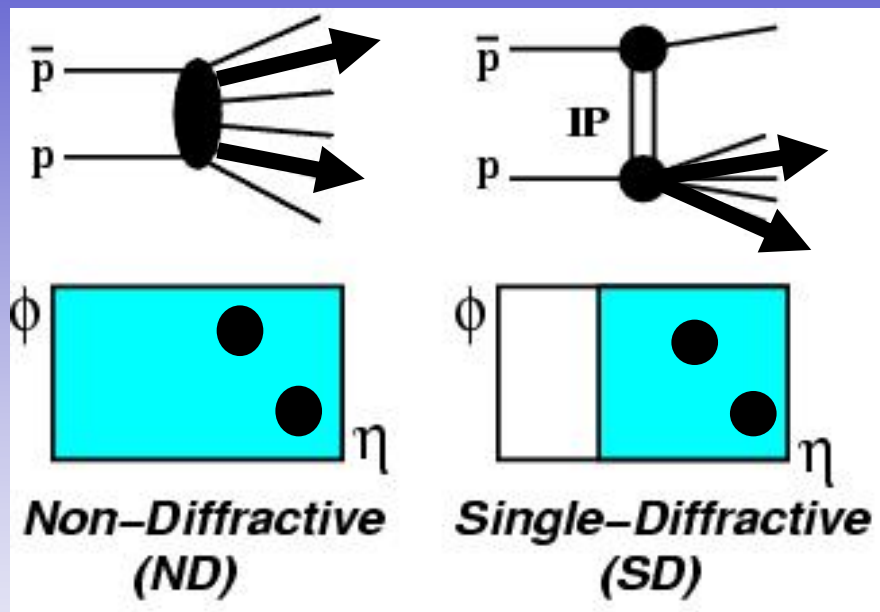


$$S = \frac{\phi \left[ \begin{array}{c} \text{yellow bar} \\ \text{white bar} \\ \text{yellow bar} \end{array} \right]_{\eta}}{\phi \left[ \begin{array}{c} \text{yellow bar} \\ \text{white bar} \\ \text{yellow bar} \end{array} \right]_{\eta}} \bigg/ \frac{\phi \left[ \text{yellow bar} \right]_{\eta}}{\phi \left[ \begin{array}{c} \text{yellow bar} \\ \text{white bar} \\ \text{yellow bar} \end{array} \right]_{\eta}}$$

$$S_{2\text{-gap}/1\text{-gap}}^{1\text{-gap}/0\text{-gap}} (1800 \text{ GeV}) \approx 0.23$$

$$S_{2\text{-gap}/1\text{-gap}}^{1\text{-gap}/0\text{-gap}} (630 \text{ GeV}) \approx 0.29$$

# DIFFRACTIVE DIJETS



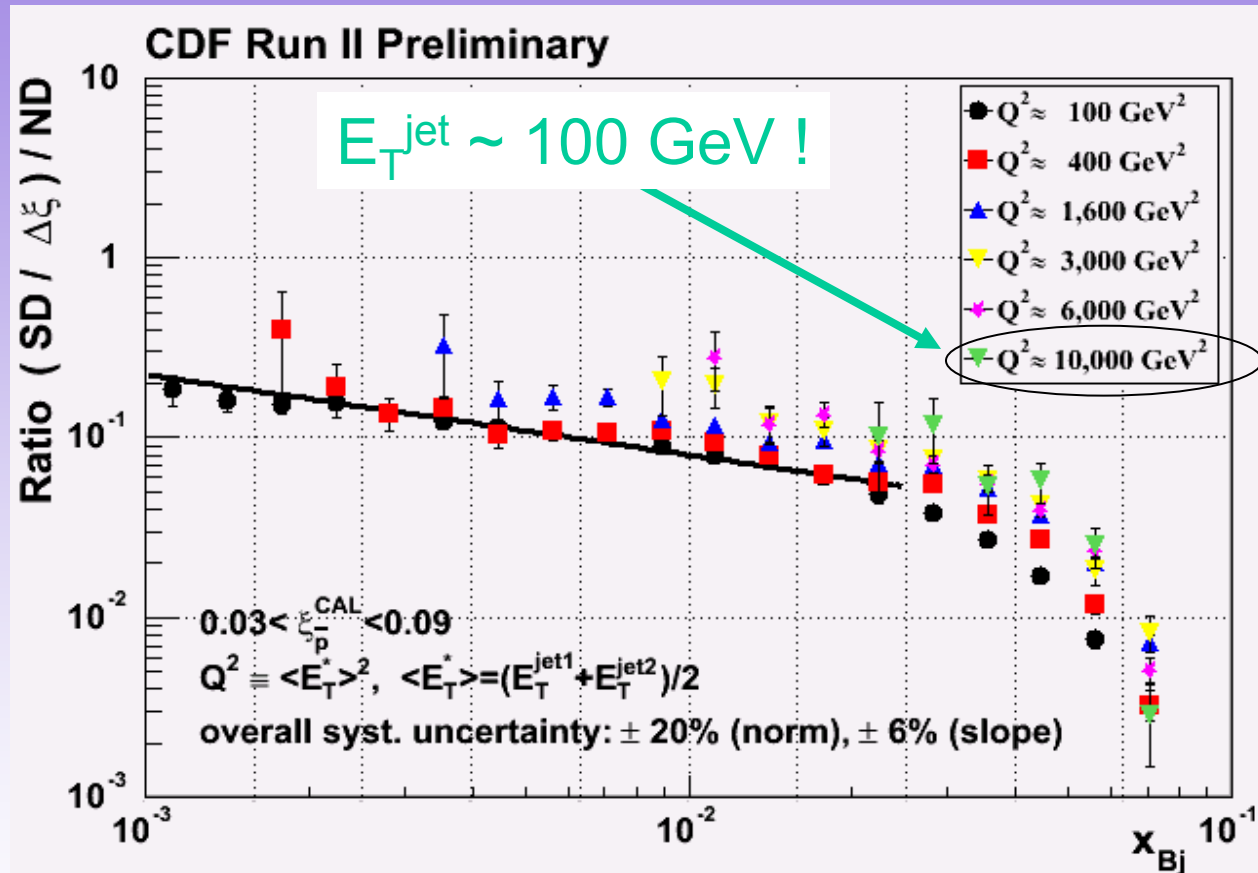
$$R(x_{Bj}) \equiv \frac{\text{Rate}_{jj}^{SD}(x_{Bj})}{\text{Rate}_{jj}^{ND}(x_{Bj})}$$

$$\Rightarrow \frac{F_{jj}^{SD}(x_{Bj})}{F_{jj}^{ND}(x_{Bj})}$$

Systematic uncertainties due to energy scale and resolution cancel out in the ratio

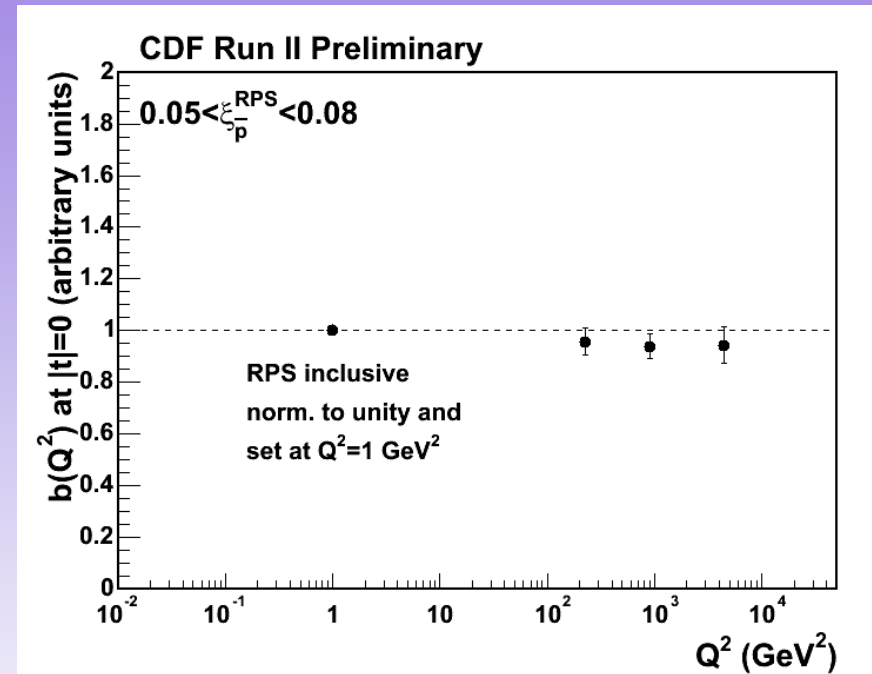
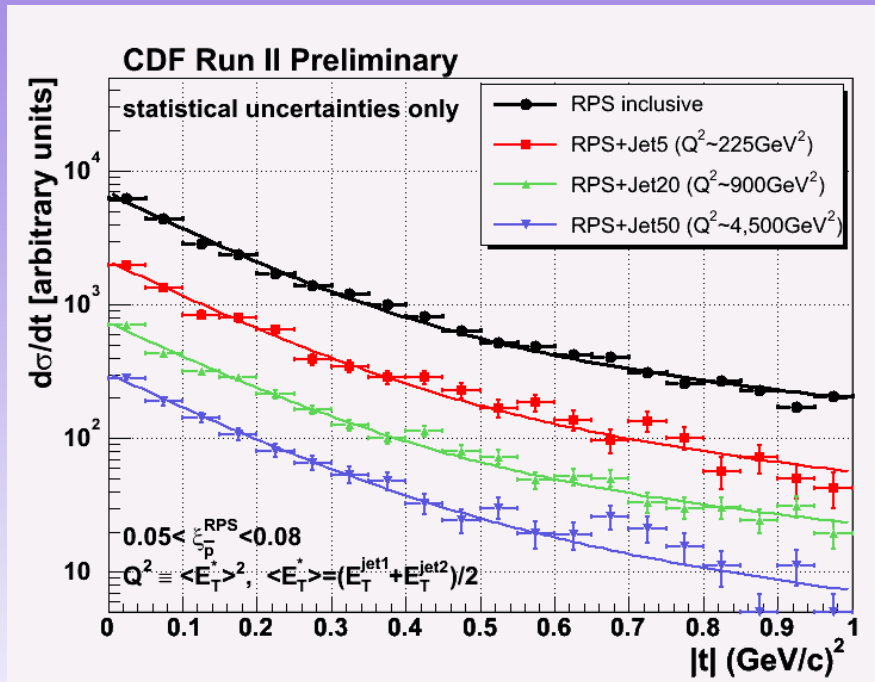
# Diffractive Structure Function:

$X_{Bj}$  and  $Q^2$  dependence



Small  $Q^2$  dependence in region  $100 < Q^2 < 10\,000 \text{ GeV}^2$   
 $\Rightarrow$  Pomeron evolves as the proton!

# Diffraction Structure Function: t- dependence



Fit  $d\sigma/dt$  to a double exponential:

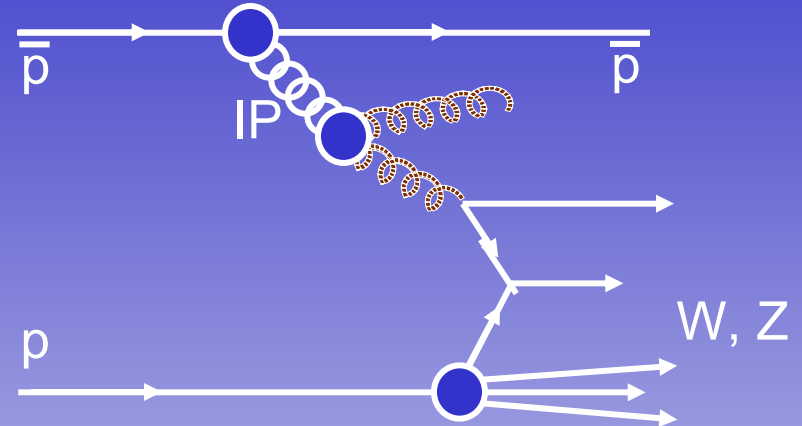
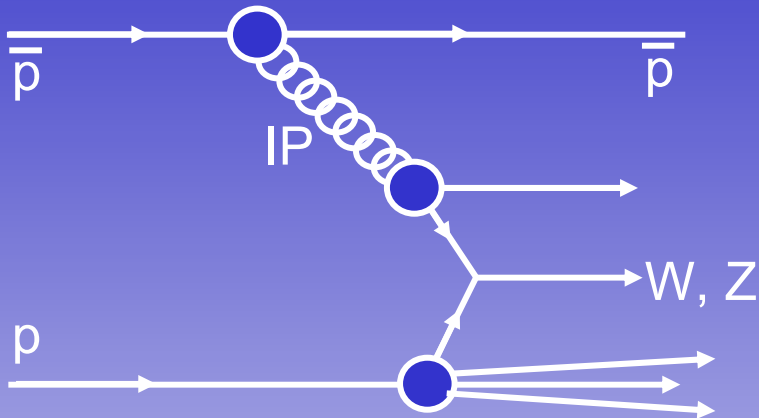
$$F = 0.9 \cdot e^{b_1 \cdot t} + 0.1 \cdot e^{b_2 \cdot t}$$

- No diffraction dips
- No  $Q^2$  dependence in slope from inclusive to  $Q^2 \sim 10^4 \text{ GeV}^2$

Remaining work:

- Obtain slope normalization
- Extend range to  $|t| \sim 4 \text{ GeV}^2$

# Diffractive W/Z production



- Diffractive W production probes the **quark content of the Pomeron**

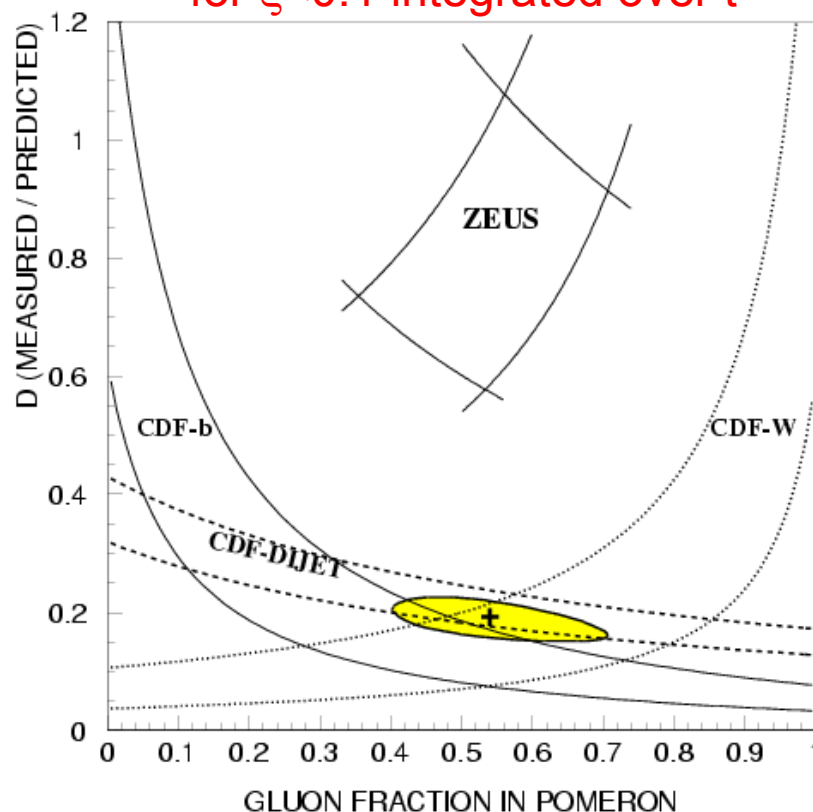
- To leading order, the W is produced by a **quark** in the Pomeron

- Production by **gluons** is **suppressed by a factor of  $\alpha_s$** , and can be distinguished from quark production by an **associated jet**

# Diffractive W/Z - motivation

- In Run I, combining diffractive dijet production with diffractive W production was used to determine the quark/gluon content of the Pomeron ==>
- In Run II, we aim at determining the diffractive structure function for a more direct comparison with HERA.
- To accomplish this we use:
  - New forward detectors
  - New methodology
  - More data

Phys Rev Lett **78**, 2698 (1997)  
Fraction of W events due to SD  
 $R^W = [1.15 \pm 0.51(\text{stat}) \pm 0.20(\text{syst})] \%$   
for  $\xi < 0.1$  integrated over  $t$





# Diffractive W/Z analysis

Using RPS information:

- ❑ No background from gaps due to multiplicity fluctuations
- ❑ No gap survival probability problem
- ❑ The RPS provides accurate event-by-event  $\xi$  measurement
- ❑ Determine the full kinematics of diffractive W production by obtaining  $\eta_v$  using the equation:

$$\xi^{\text{RPS}} - \xi^{\text{cal}} = \frac{E_T}{\sqrt{s}} e^{-\eta_v} \quad \text{where} \quad \xi^{\text{cal}} = \sum_{\text{towers}} \frac{E_T}{\sqrt{s}} e^{-\eta}$$

This allows determination of:

- W mass
- $X_{Bj}$
- Diffractive structure function

# W/Z selection requirements

## Standard W/Z selection

$$E_T^e(p_T^\mu) > 25 \text{ GeV}$$

$$\cancel{E}_T > 25 \text{ GeV}$$

$$40 < M_T^W < 120 \text{ GeV}$$

$$|Z_{\text{vtx}}| < 60 \text{ cm}$$

$$E_T^{e1}(p_T^{\mu1}) > 25 \text{ GeV}$$

$$E_T^{e2}(p_T^{\mu2}) > 25 \text{ GeV}$$

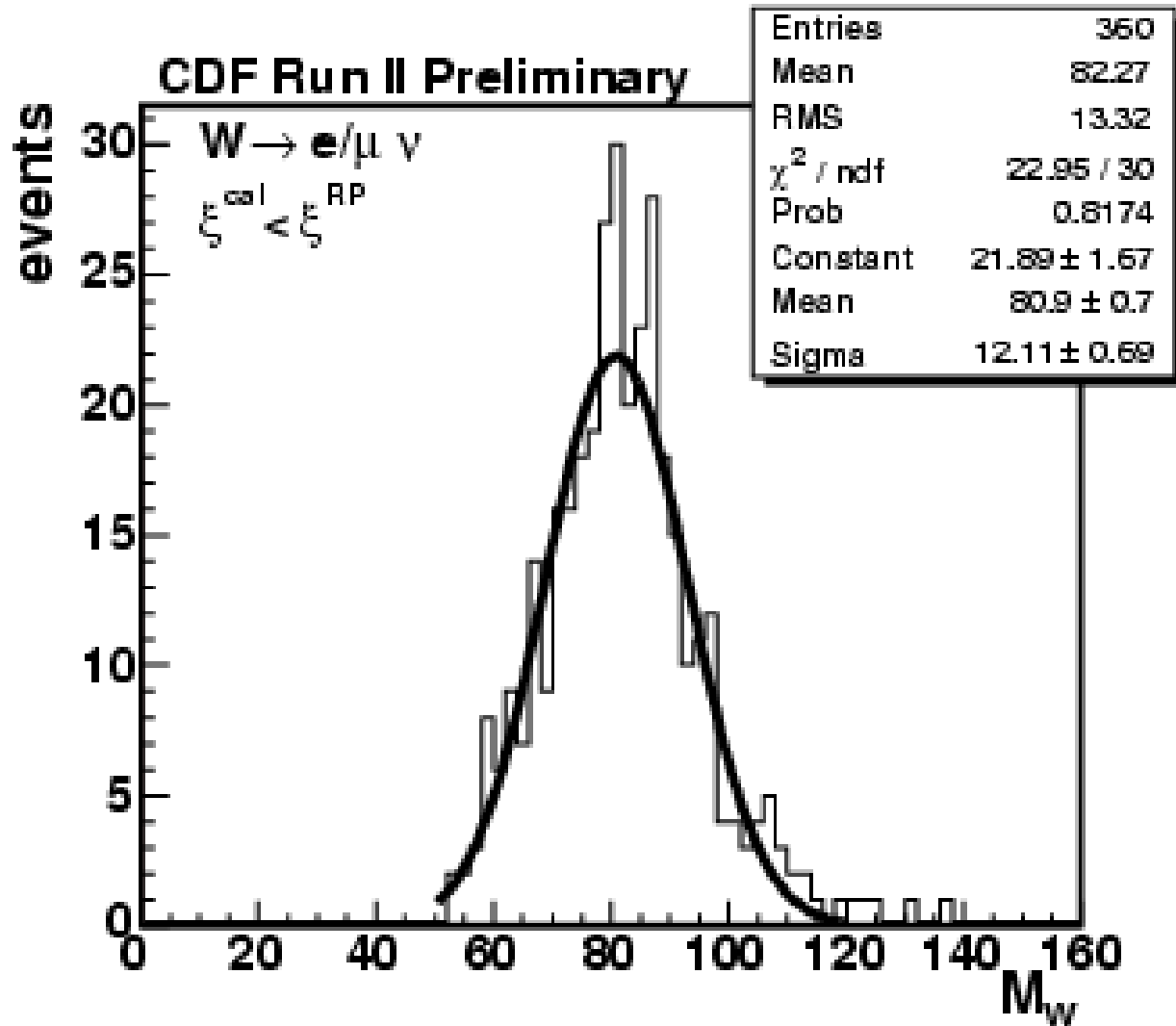
$$66 < M^Z < 116 \text{ GeV}$$

$$|Z_{\text{vtx}}| < 60 \text{ cm}$$

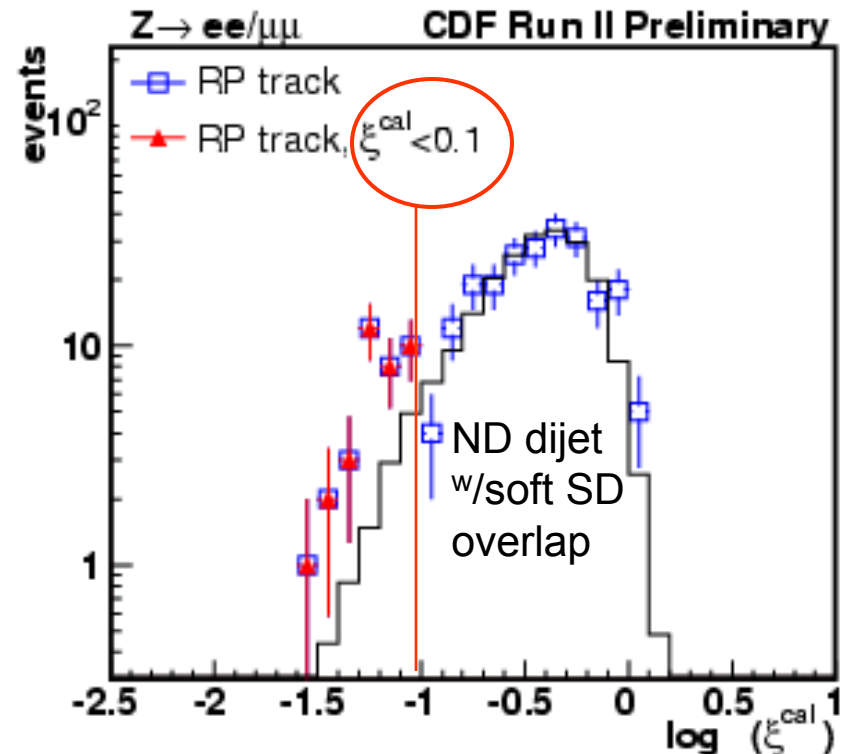
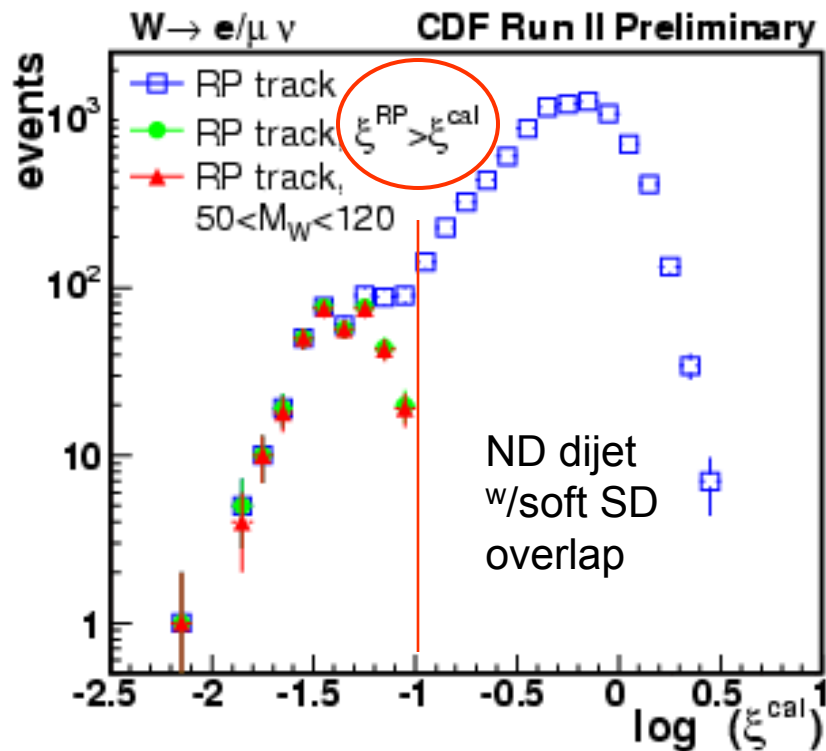
## Diffraction W/Z selection

- RPS trigger counters - MIP
- RPS track -  $0.03 < \xi < 0.10$ ,  $|t| < 1$
- $W \rightarrow 50 < M_W(\xi^{\text{RPS}}, \xi^{\text{cal}}) < 120$
- $Z \rightarrow \xi^{\text{cal}} < 0.1$

# Reconstructed Diffractive W-Mass



# Rejection of Multiple Interactions



# Diffractive W/Z results

$$R^W (0.03 < \xi < 0.10, |t| < 1) = [0.97 \pm 0.05(\text{stat}) \pm 0.11(\text{syst})]\%$$

Run I:  $R^W = 1.15 \pm 0.55\%$  for  $\xi < 0.1 \rightarrow$  estimate  $0.97 \pm 0.47\%$  in  $0.03 < \xi < 0.10$  &  $|t| < 1$ )

$$R^Z (0.03 < x < 0.10, |t| < 1) = [0.85 \pm 0.20(\text{stat}) \pm 0.11(\text{syst})]\%$$

## *CDF/DØ Comparison – Run I ( $\xi < 0.1$ )*

CDF PRL 78, 2698 (1997)

$$R^W = [1.15 \pm 0.51(\text{stat}) \pm 0.20(\text{syst})]\%$$

gap acceptance  $A^{\text{gap}} = 0.81$

uncorrected for  $A^{\text{gap}} \rightarrow$

$$R^W = (0.93 \pm 0.44)\%$$

( $A^{\text{gap}}$  calculated from MC)

DØ Phys Lett B 574, 169 (2003)

$$R^W = [5.1 \pm 0.51(\text{stat}) \pm 0.20(\text{syst})]\%$$

gap acceptance  $A^{\text{gap}} = (0.21 \pm 4)\%$

uncorrected for  $A^{\text{gap}} \rightarrow$

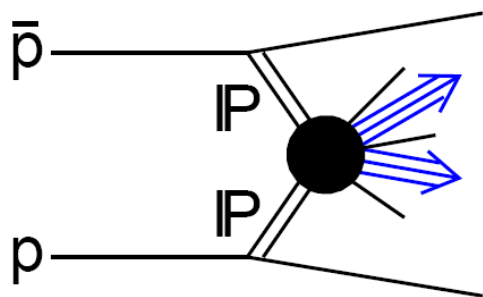
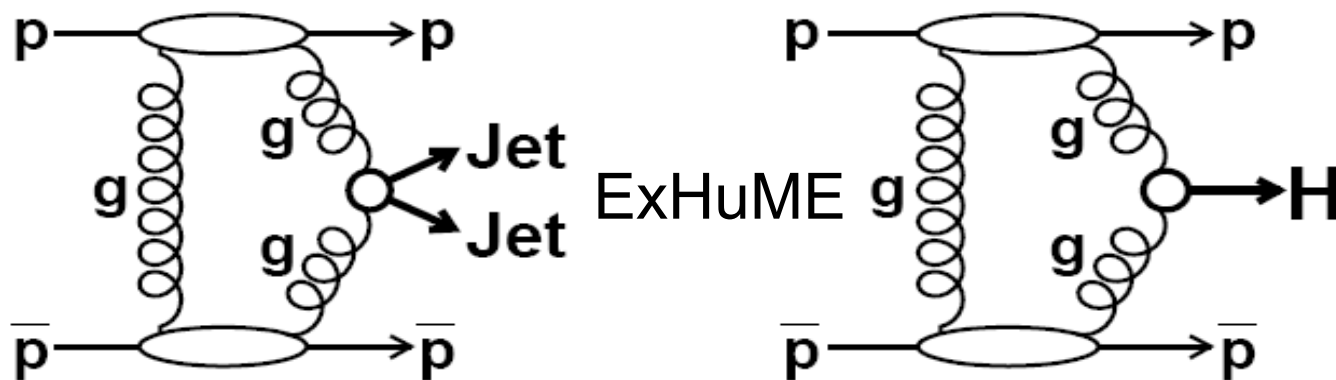
$$R^W = [0.89 + 0.19 - 0.17]\%$$

$$R^Z = [1.44 + 0.61 - 0.52]\%$$

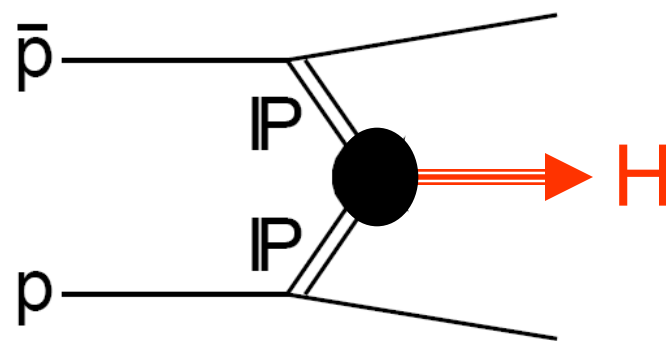
Stay connected for  $F^D_{W/Z}$

# Exclusive Dijet and Higgs Production

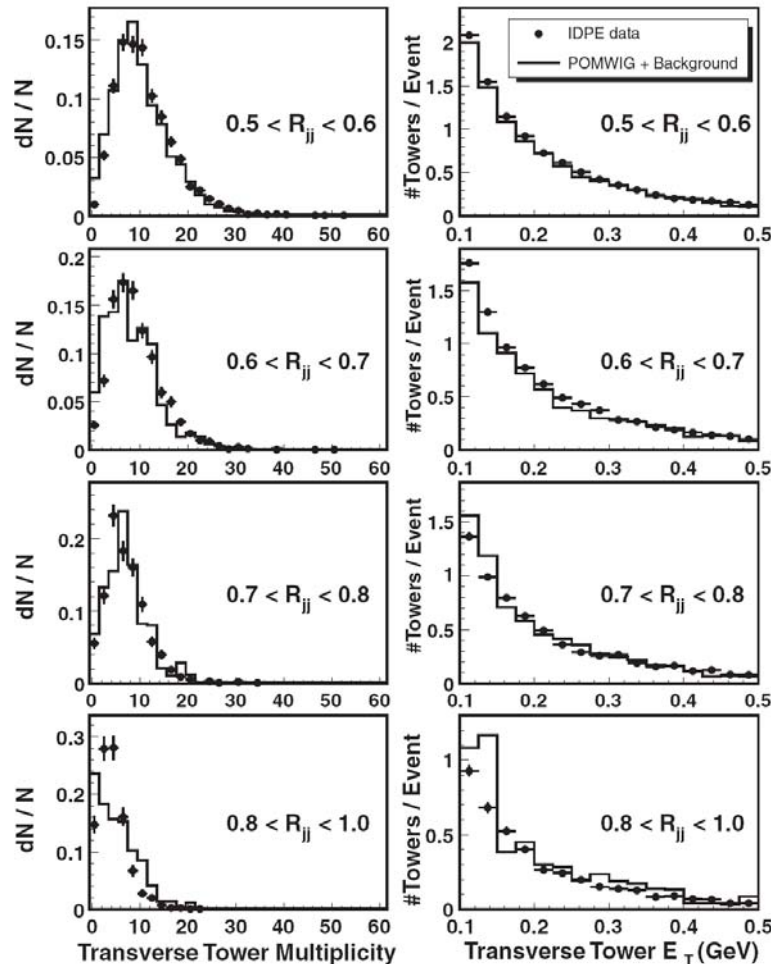
Phys. Rev. D 77, 052004



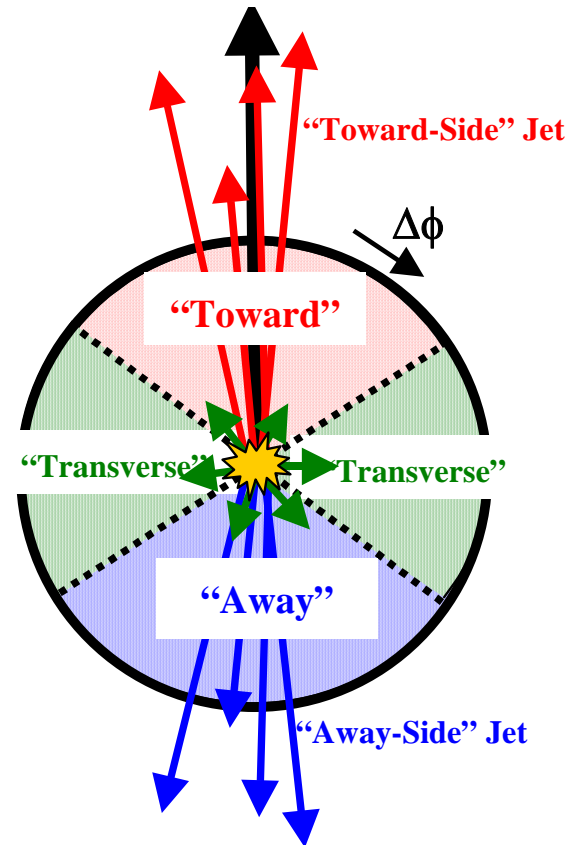
DPEMC



# Underlying Event (UE)



Is it modeled correctly?



The data and POMWIG+Background distributions in the transverse  $\Delta\phi$ -region relative to the di-jet axis agree, indicating that the UE is correctly modeled.

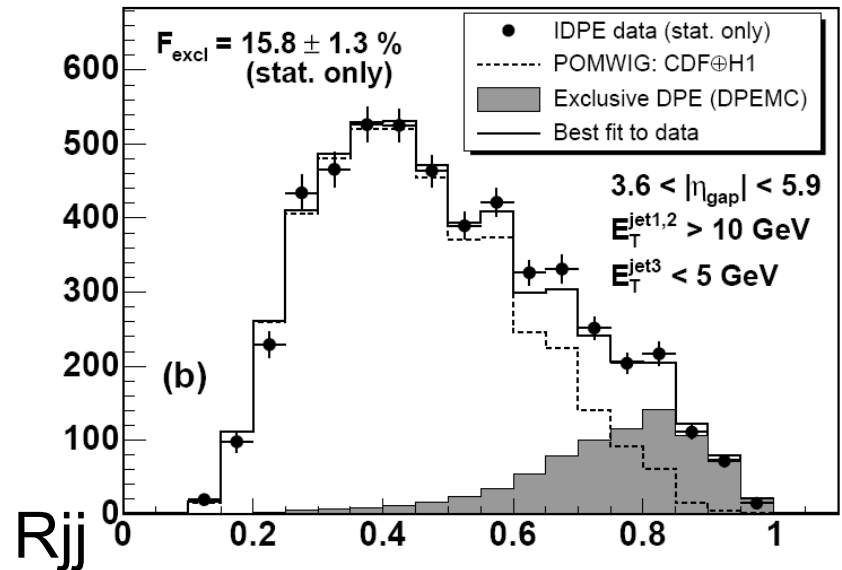
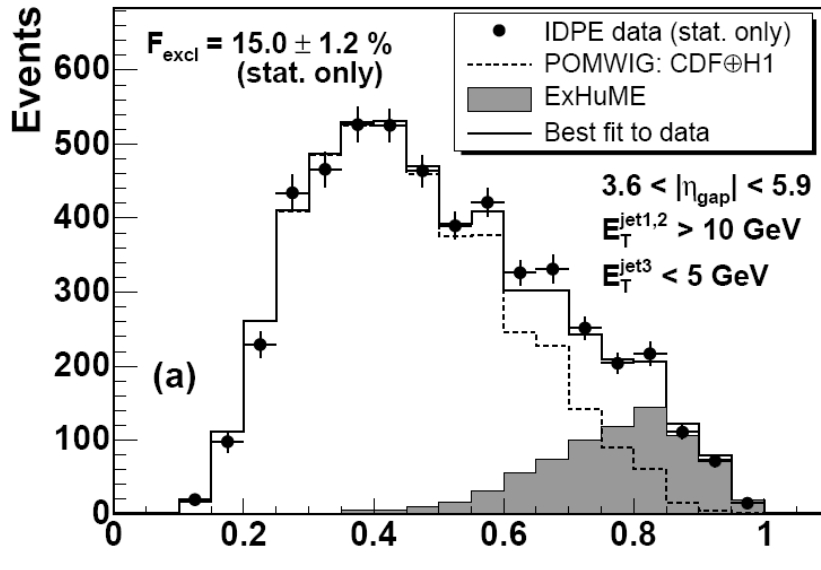


# Inclusive DPE $W/LRG_p$ : Data vs. MC

ExHuME

←exclusive MC models→

DPEMC

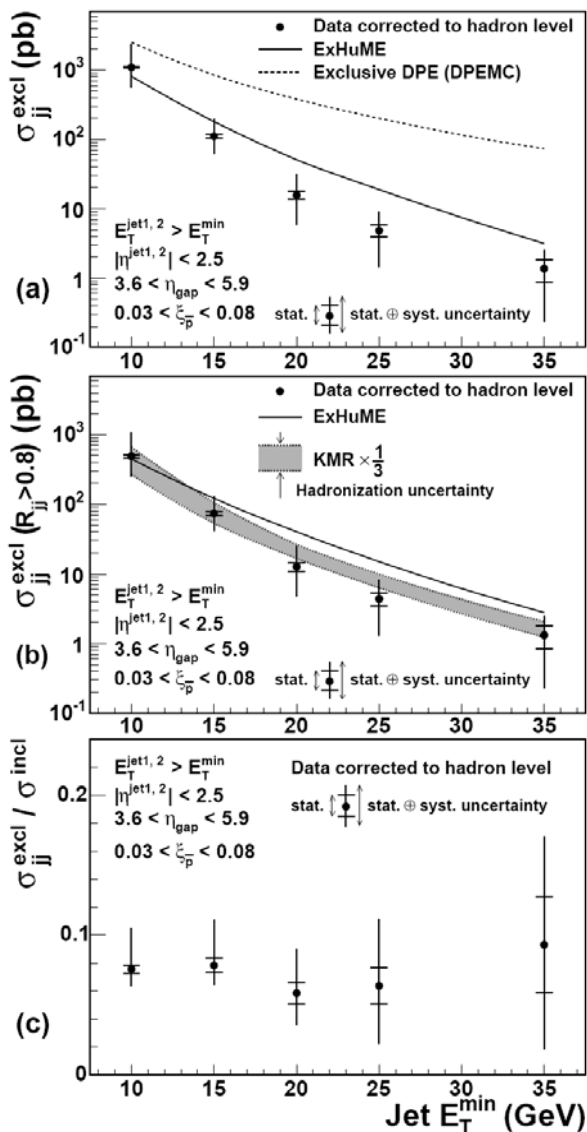


ExHuME (KMR):  $gg \rightarrow gg$  process  
(based on LO pQCD)

DPEMC: exclusive DPE MC  
based on Regge theory

Shape of excess of events at high  $R_{jj}$   
is well described by both ExHuME & DPEMC – but...

# ExHuME vs. DPEMC and vs. data



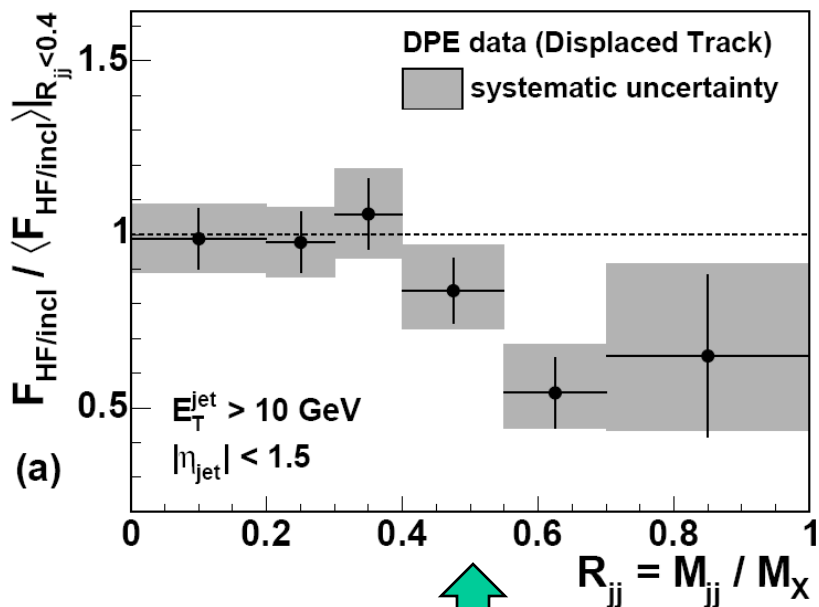
Measured x-sections favor ExHuME

KMR  $\times 1/3$  agrees with data  
 → Within theoretical uncertainty of +/- factor of 3

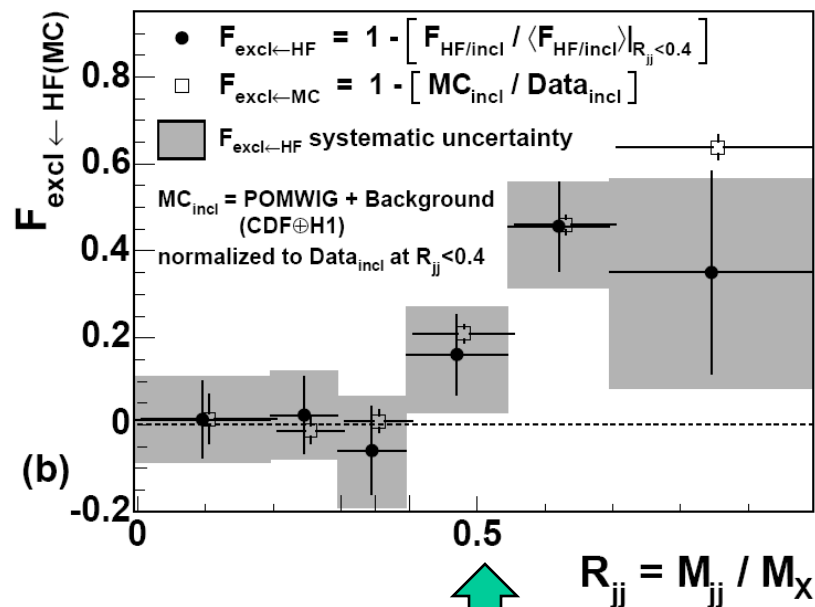
$\sigma_{jj}^{\text{excl}} / \sigma_{jj}^{\text{incl}}$  approx. independent of  $E_T^{\text{min}}$

# Heavy Flavor suppression vs. Inclusive Signal

## HF suppression

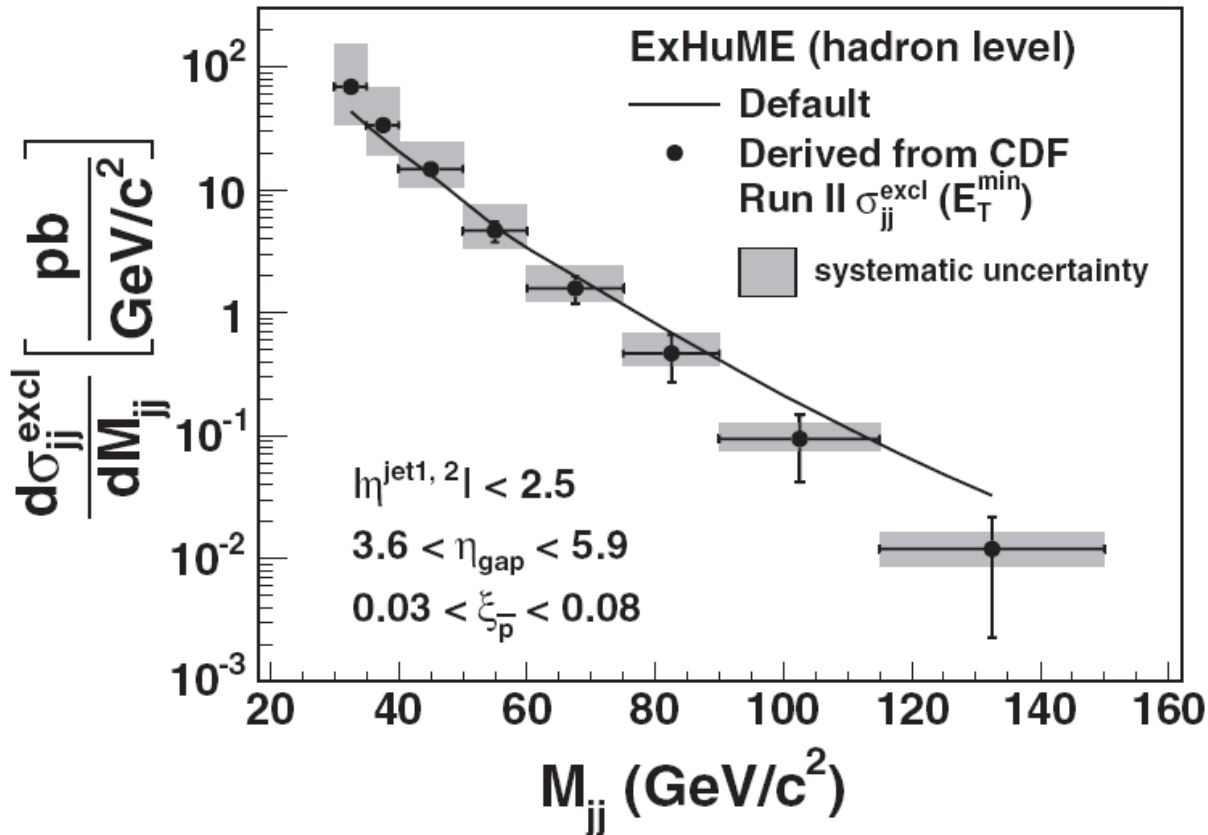


## HF suppression vs. Incl



Invert HF vertically and compare with 1-MC/DATA  
 → good agreement observed

# Exclusive Dijet x-section vs. $M_{jj}$



line: ExHuME hadron-level exclusive di-jet cross section vs. di-jet mass  
points: derived from CDF excl. di-jet x-sections using ExHuME

Stat. and syst. errors are propagated from measured cross section uncertainties using  $M_{jj}$  distribution shapes of ExHuME generated data.

# Exclusive $\ell$ and $\gamma\gamma$

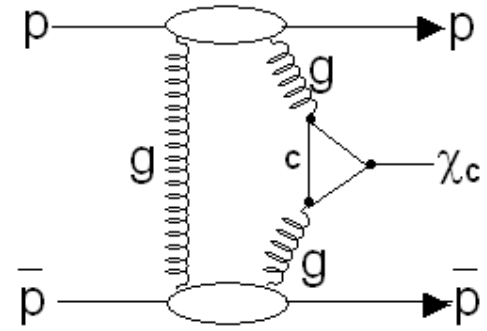
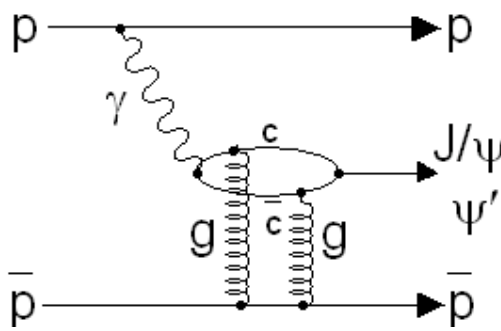
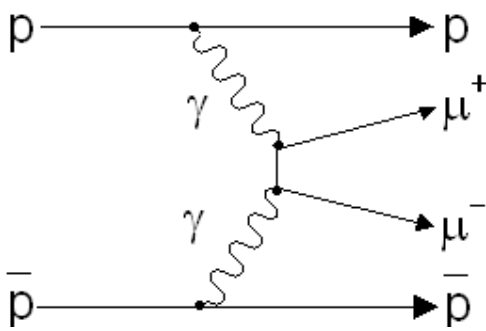
$e^+e^- (> 10)$ ,  $\gamma\gamma (> 10)$ ,  $\mu^+\mu^- (3-4)$ ,  $J/\psi, \psi(2S), \chi_c^0, l^+l^- (> 40)$ ,  $Z$

PRL98, 112001  
2007

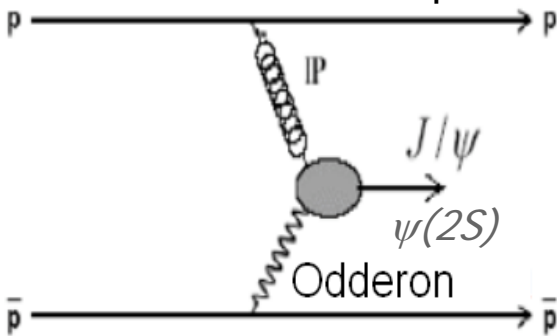
PRL99, 242002  
2007

PRL under coll. review

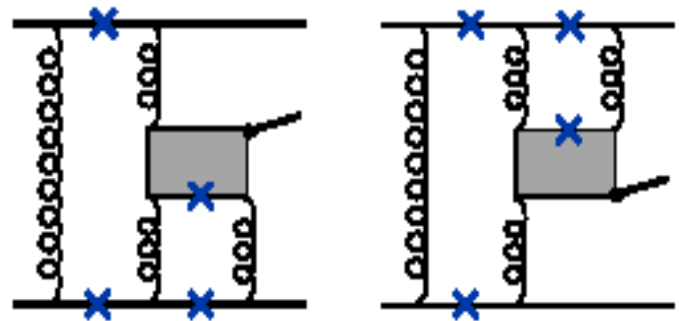
PRL under coll. review



**Odderon:** C-odd partner of Pomeron  $\rightarrow$  3 gluons in a color-neutral state



LO diagrams  $\rightarrow$



# Exclusive $J/\psi$ and $\psi(2s)$ production\*

## $J/\psi$ production

$$d\sigma/dy|_{y=0} = 3.92 \pm 0.62 \text{ nb}$$

*In agreement with av. Prediction of  $3.0 \pm 0.3$  nb*

2.8 nb [Szczyrek07,], 2.7 nb [Klein&Nystrand04],  
3.0 nb [Conclaves&Machado05], and  
3.4 nb [Motkya&Watt08].

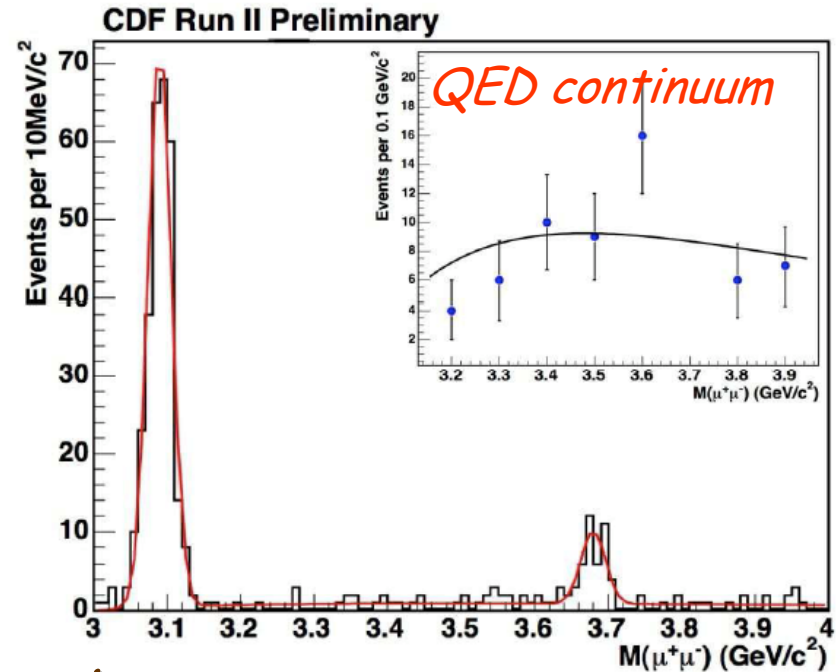
## $\Psi(2s)$ production

$$d\sigma/dy|_{y=0} = 0.54 \pm 0.15 \text{ nb}$$

$$R = \psi(2s)/J/\psi = 0.14 \pm 0.05$$

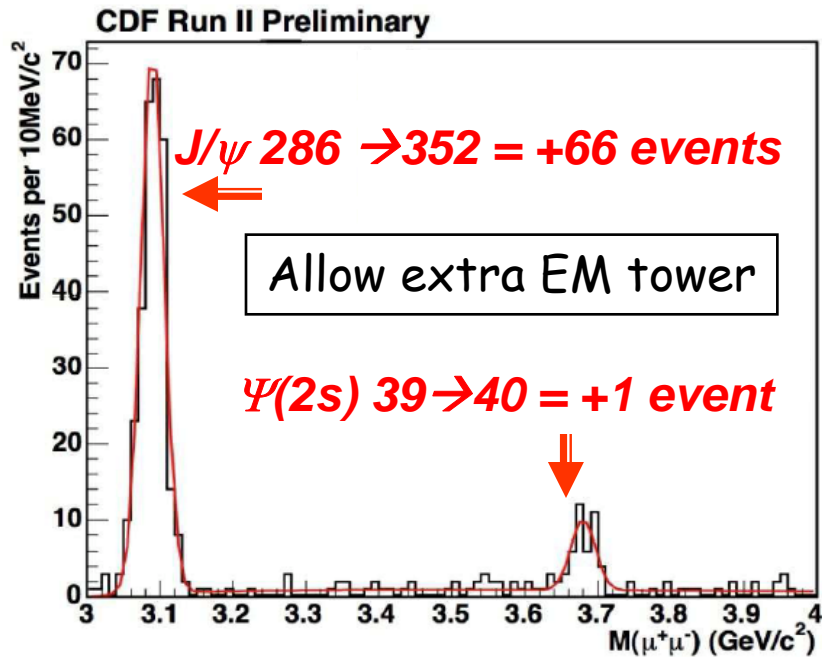
*In agreement with HERA:*

*$R = 0.166 \pm 0.012$  in a similar kinematic region*



\* James Pinfold <http://www.fp420.com/conference/dec2008/index.html>

# Exclusive $\chi_c \rightarrow J/\psi(\mu^+\mu^-) + \gamma$ production\*



- *Allowing EM towers (with  $EmE_T > 80$  MeV gives a large increase in the  $J/\psi$  peak but a minor change in the  $\psi(2s)$  peak*
- $\Rightarrow$  *Evidence for  $\chi_c \rightarrow J/\psi + \gamma$  production*

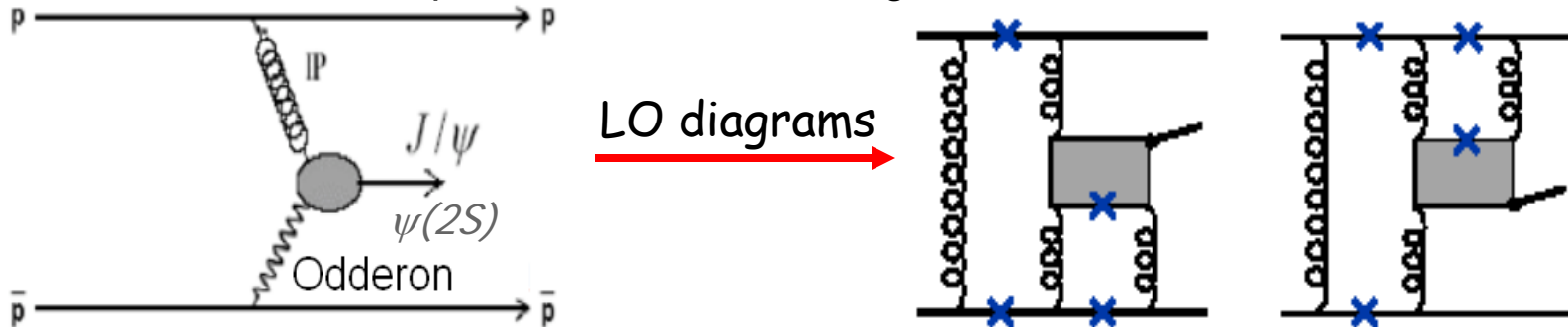
$d\sigma/dy|_{y=0} = 75 \pm 14$  nb, compatible with theoretical predictions of 150 nb (Yuan 01) & 130 nb (KRS01) - error of  $\mathcal{O}(50)$  nb

\* James Pinfold <http://www.fp420.com/conference/dec2008/index.html>



# Exclusive $\mu^+\mu^-$ and the Odderon\*

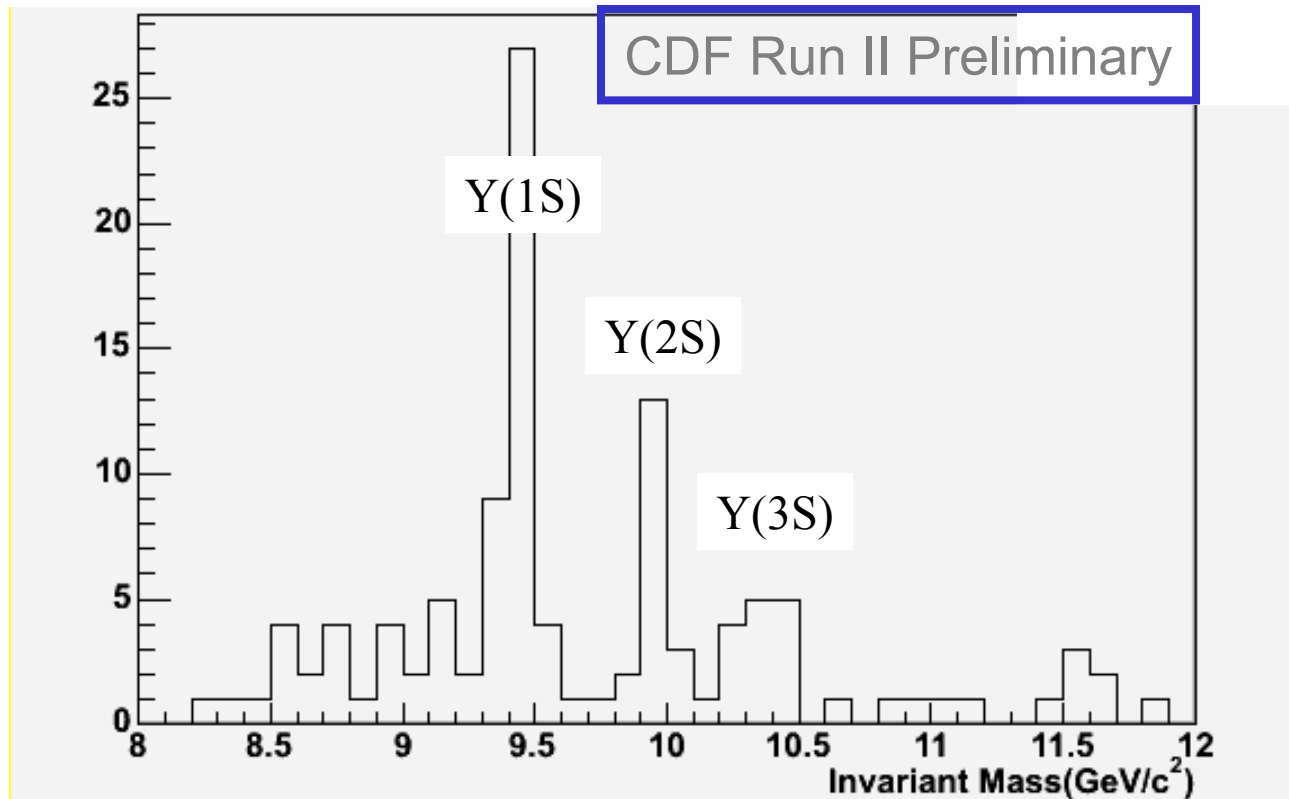
**Odderon:** C-odd partner of Pomeron  $\rightarrow$  3 gluons in a color-neutral state



- ❑ *The odderon would contribute to  $J/\psi$ ,  $\psi(2s)$  (&  $Y$ ) peaks -not the  $\chi_c$*
- ❑ *The  $J/\psi$  &  $\psi(2s)$  cross-sections agree with predictions (that fit the HERA data)  $\rightarrow$  no significant odderon signa*
  - ❑  *$R(\text{exp./theory})_{J/\psi} = 1.32 \pm 0.41$  ,  $R(\text{exp./theory})_{\psi(2s)} = 1.15 \pm 0.21$*
  - $\rightarrow R(\text{data/theory}) [\text{combined } J/\psi \text{ \& } \psi(2s)] = 1.19 \pm 0.19$*
- ❑ *Limit on odderon prod. -  $R[(O\text{-}IP) \rightarrow V / (\gamma\text{-}IP) \rightarrow V] < 0.34$  (95% CL)*
- ❑ *Another limit:  $R[(O+\gamma)IP \rightarrow J/\psi / IP\text{-}IP \rightarrow \chi_c(3415)] < 0.060 \pm 0.015$*

\* James Pinfold <http://www.fp420.com/conference/dec2008/index.html>

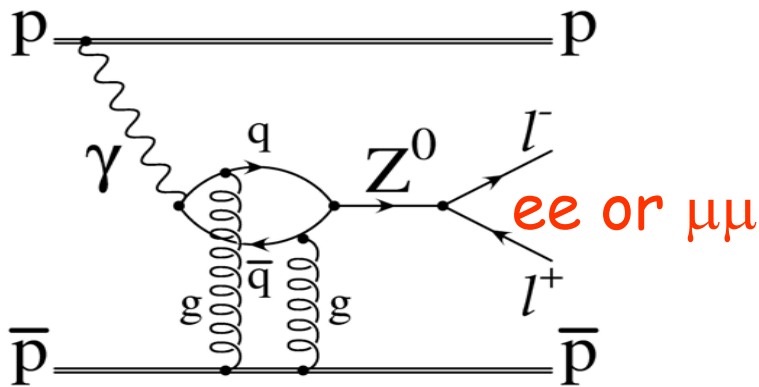
# Exclusive Upsilon photoproduction\*



$M(\mu^+\mu^-)$  with  $p_T(\mu^+\mu^-) < 1.5$  GeV/c and  $\pi - \Delta\phi < 0.34$  rad

\* Michael Albrow <http://www.fp420.com/conference/dec2008/index.html>

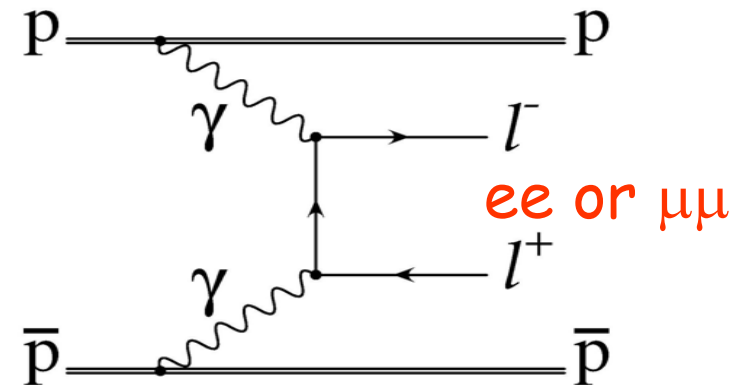
# Exclusive $Z / \gamma\gamma \rightarrow \ell\ell$ studies\*



<sup>1</sup> Phys. Rev. D78, 014023 (2008)

<sup>2</sup> Eur. Phys. J., C56 33 (2008)

<sup>3</sup> Phys. Rev. D72,036007



Exclusive Z SM cross section:

0.3 fb [Motyka & Watt]<sup>1</sup>

1.3 fb [Goncalves & Machado]<sup>2</sup>

Expect  $\sim 0.6 - 2.6$  events in 2 fb<sup>-1</sup>  
(not including 3.37% leptonic BF).

→ Search for BSM physics,  
e.g. color sextet quark model<sup>3</sup>: much  
enhanced cross section expected.

Exclusive di-lepton cross section:

➤ Background to exclusive Z

➤ Can be used to calibrate forward  
proton detectors:

$$\xi = s^{-1/2} \sum p_T |e^{-\eta}|$$

\* Emily Nurse <http://www.fp420.com/conference/dec2008/index.html>

# Exclusive $Z / \gamma\gamma \rightarrow \ell\ell$ results

$\mathcal{L} = 2.20$  (2.03)  $\text{fb}^{-1}$  in the electron (muon) channels

$e+e^-$  or  $\mu+\mu^-$  with  $p^T > 25$   
 $\text{GeV}$

$|\eta_{\mu}^1| < 1.0, |\eta_{\mu}^2| < 1.5$

$|\eta_e^1| < 1.3, |\eta_e^2| < 3.6$

require  $82 < M_{\ell\ell} < 98 \text{ GeV}$

$W \rightarrow \ell\nu$  events used as a control to study exclusivity cuts

CDF Run II Preliminary

## Exclusive di-leptons

$$\sigma(pp \rightarrow p \ell\ell p) = 0.24^{+0.13}_{-0.10} \text{ pb}$$

$M_{\ell\ell} > 40 \text{ GeV}, |\eta_d| < 4.0$

LPAIR prediction:  $\sigma = 0.256 \text{ pb}$

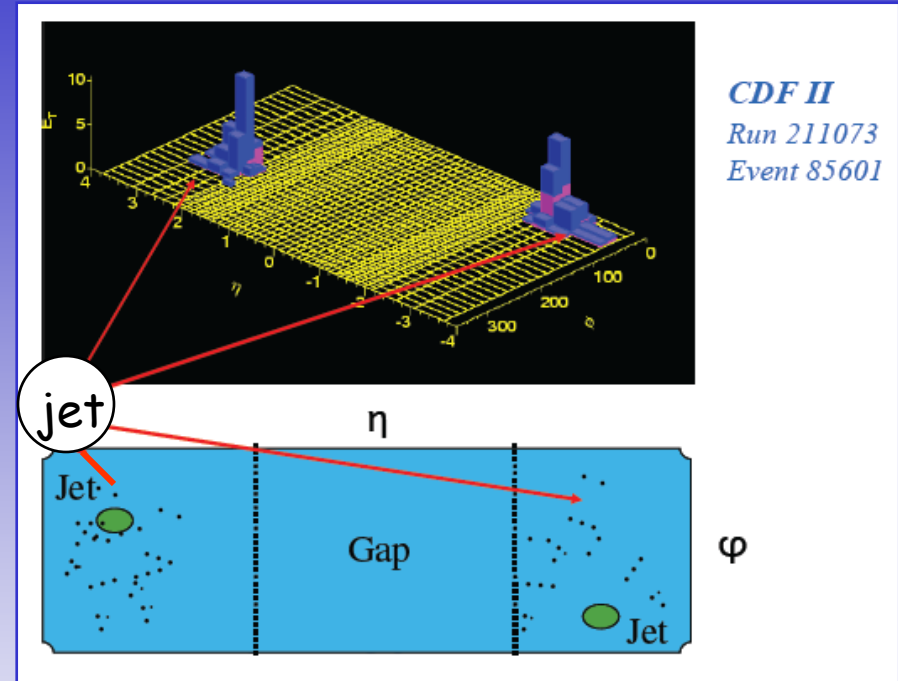
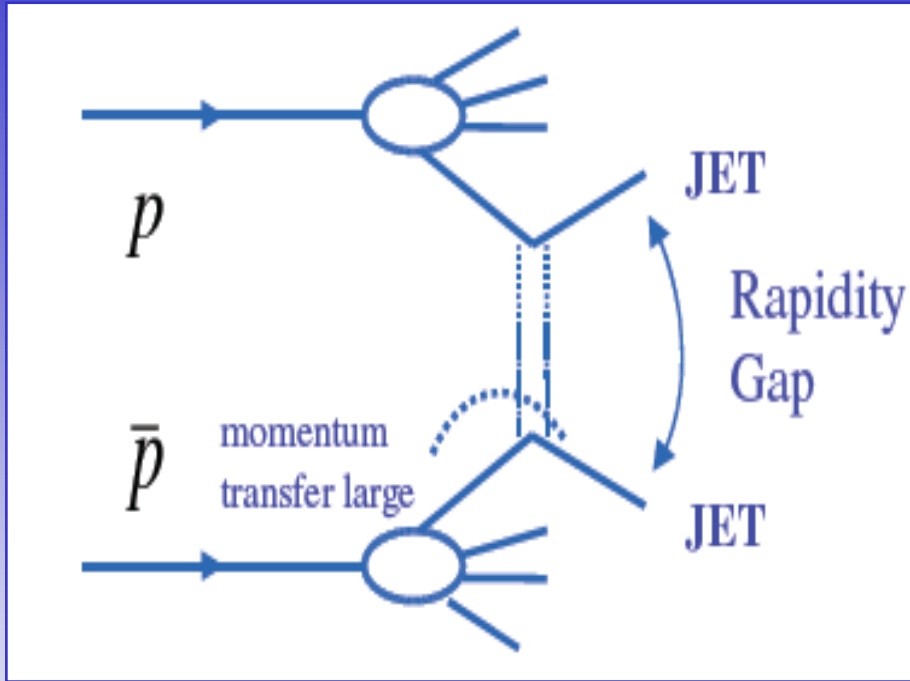
## Exclusive Z

$$\sigma(Z_{\text{excl}}) < 0.96 \text{ pb @ 95\% C.L.}$$

SM: 0.3 fb [Motyka & Watt]

1.3 fb [Goncalves & Machado]

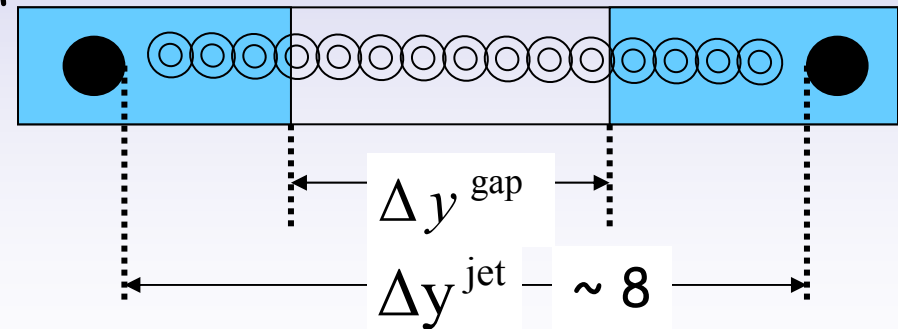
# CENTRAL RAPIDITY GAPS \*



- Measure  $\Delta y^{\text{gap}}$  width and position to differentiate among models.

$$\Delta y^{\text{gap}} = \Delta y^{\text{jet}} \Rightarrow \text{BFKL}$$

$$\Delta y^{\text{gap}} < \Delta y^{\text{jet}} \Rightarrow \text{composite}$$



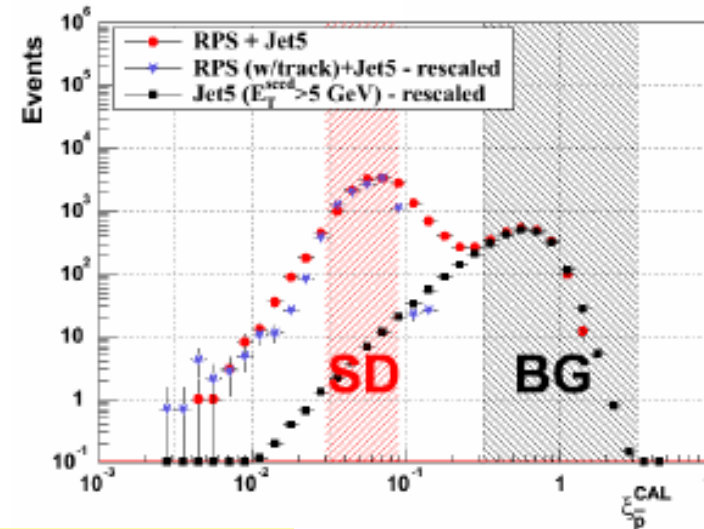
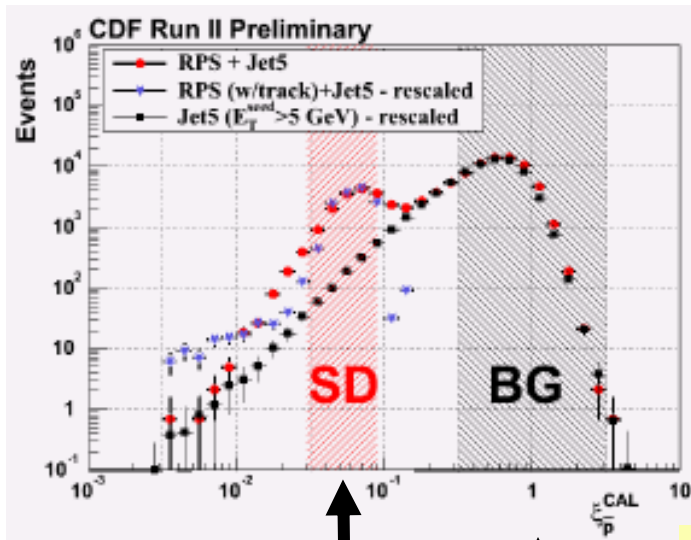
\* Christina Mesropian <http://www.cs.infn.it/diff2008/program.html>

# Low Luminosity Run

→ January 2006: data with dedicated diffractive triggers ←

2002-303 data  $\sim 1.5E31$

Low Lum  $\sim 0.5E30$



single diffraction  
 $0.03 < \xi < 0.1$

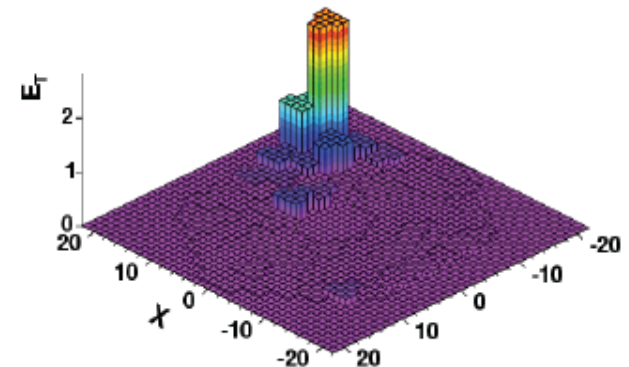
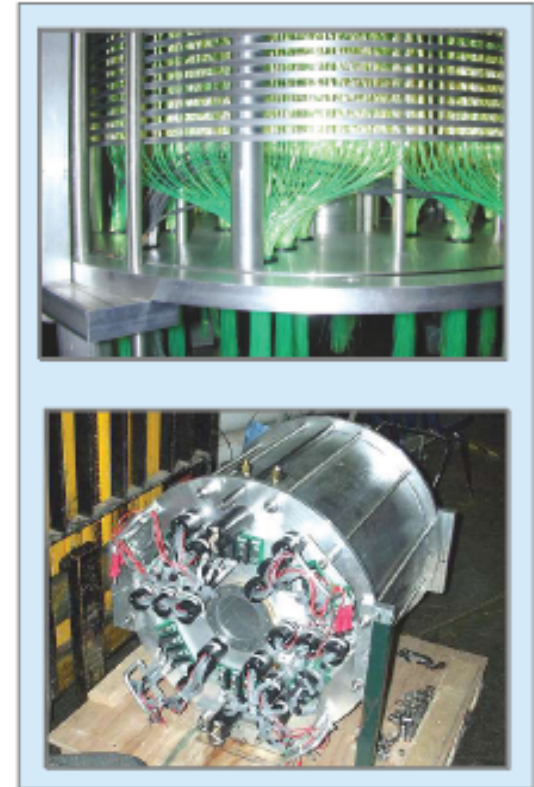
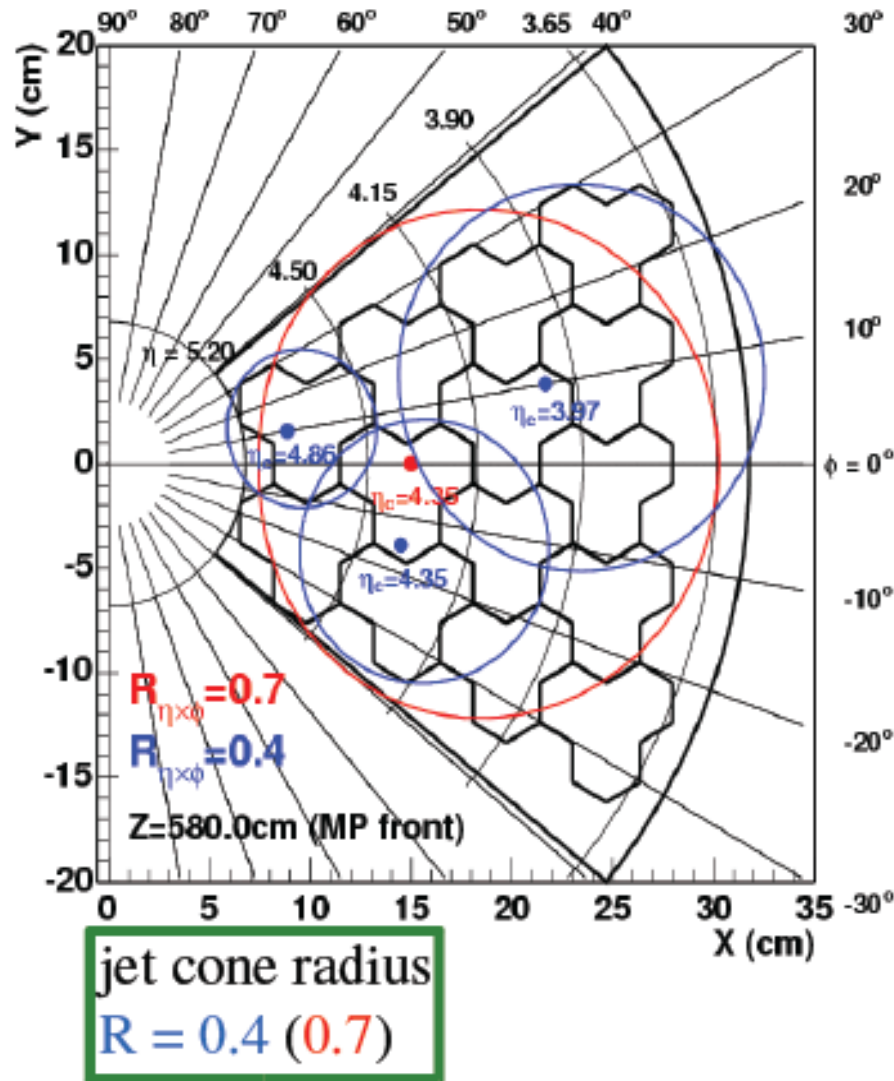
$$\xi^{\text{cal}} = \sum_{\text{towers}} \frac{E_T}{\sqrt{s}} e^{-\eta}$$

diffractive dijet overlapped with MB soft diffraction



# MiniPlug Jets

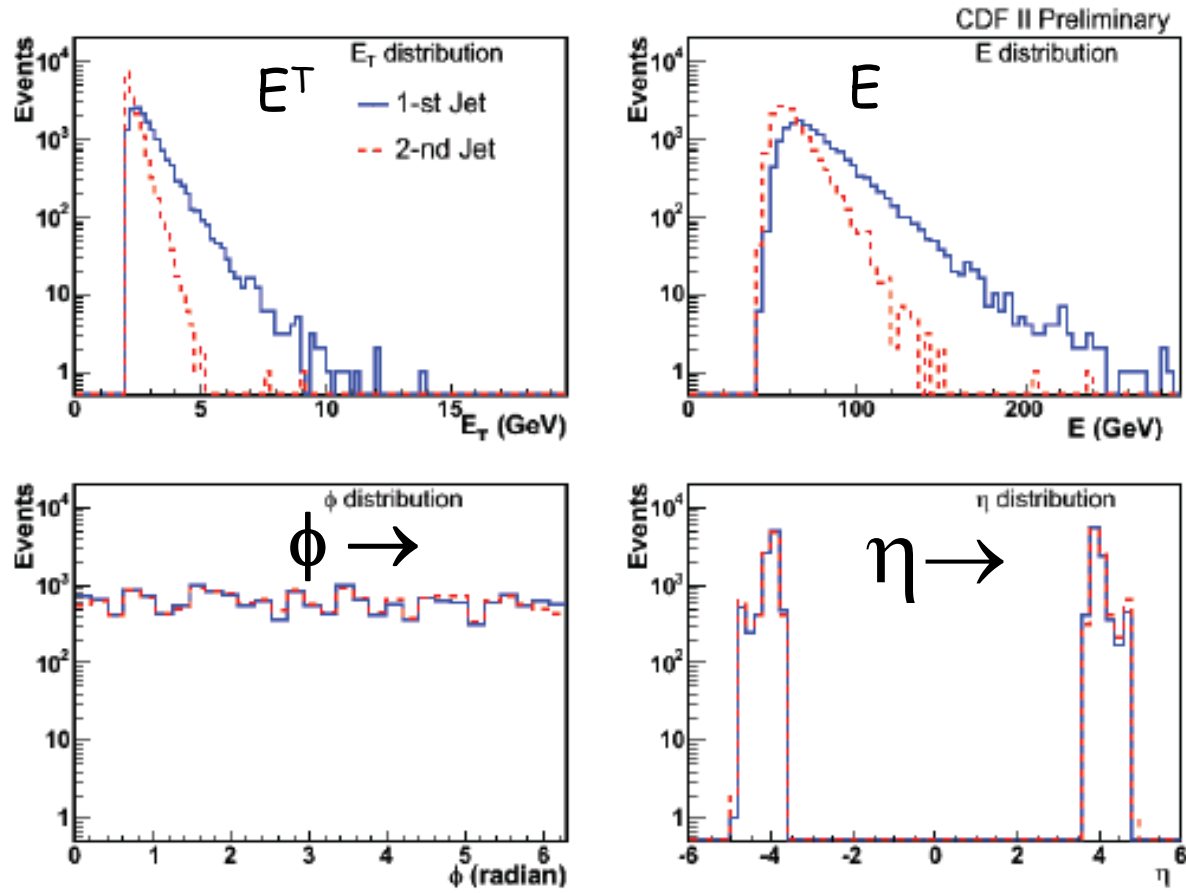
*Nucl. Instrum. Meth. A518 (2004) 42.*  
*Nucl. Instrum. Meth. A496 (2003) 333.*



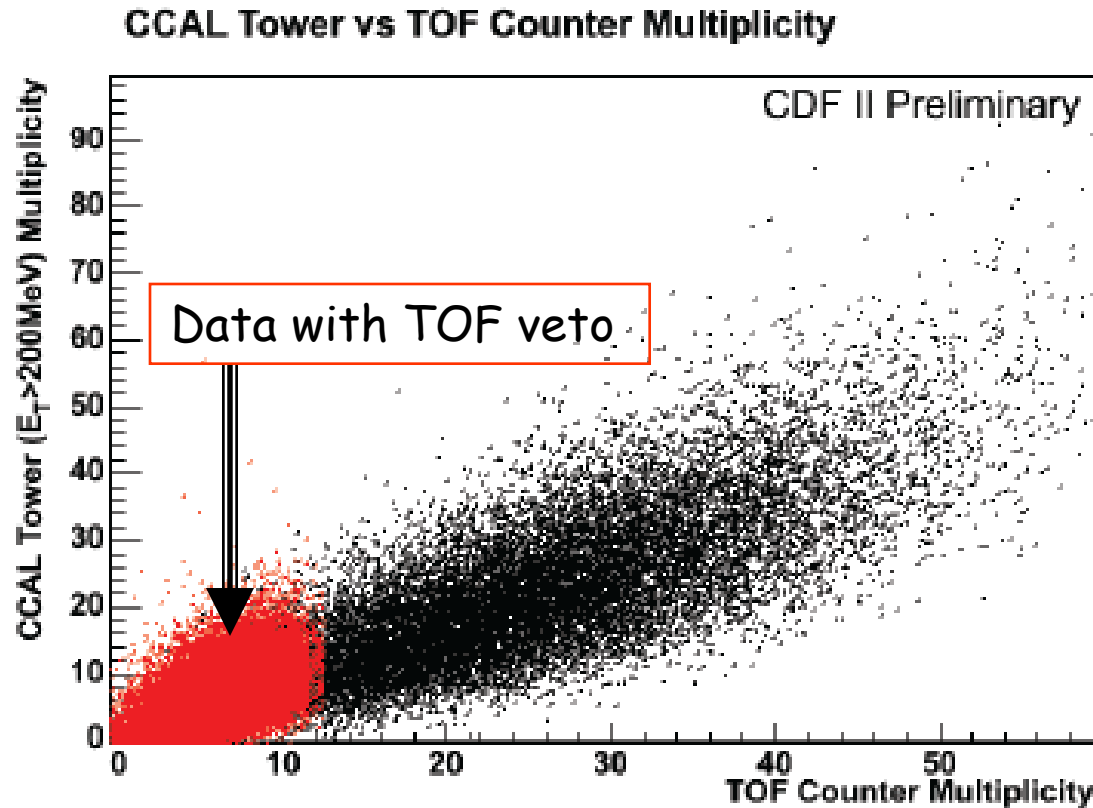


# MiniPlug Jet Properties

$$E_{T}^{\text{jet}1,2} > 2 \text{ GeV}, 3.5 < |\eta|^{\text{jet}1,2} < 5.1, \eta^{\text{jet}1} \cdot \eta^{\text{jet}2} < 0$$

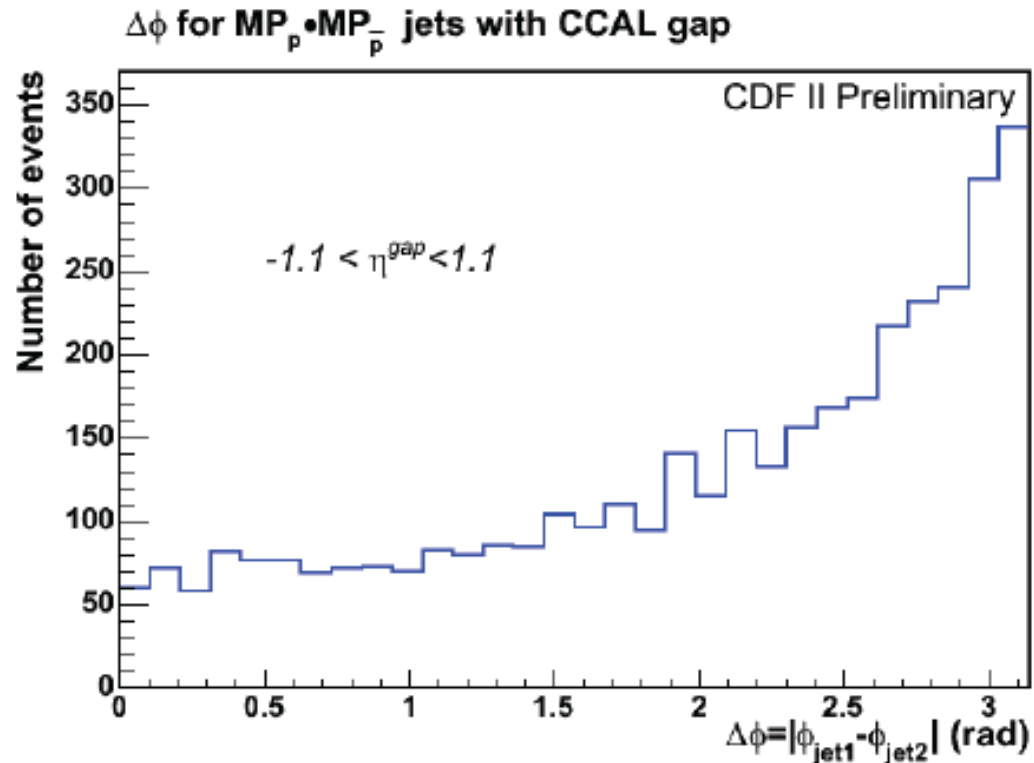
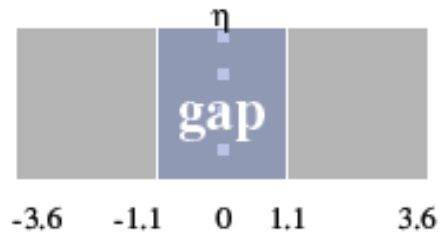


# MP Jet Data with TOF Veto



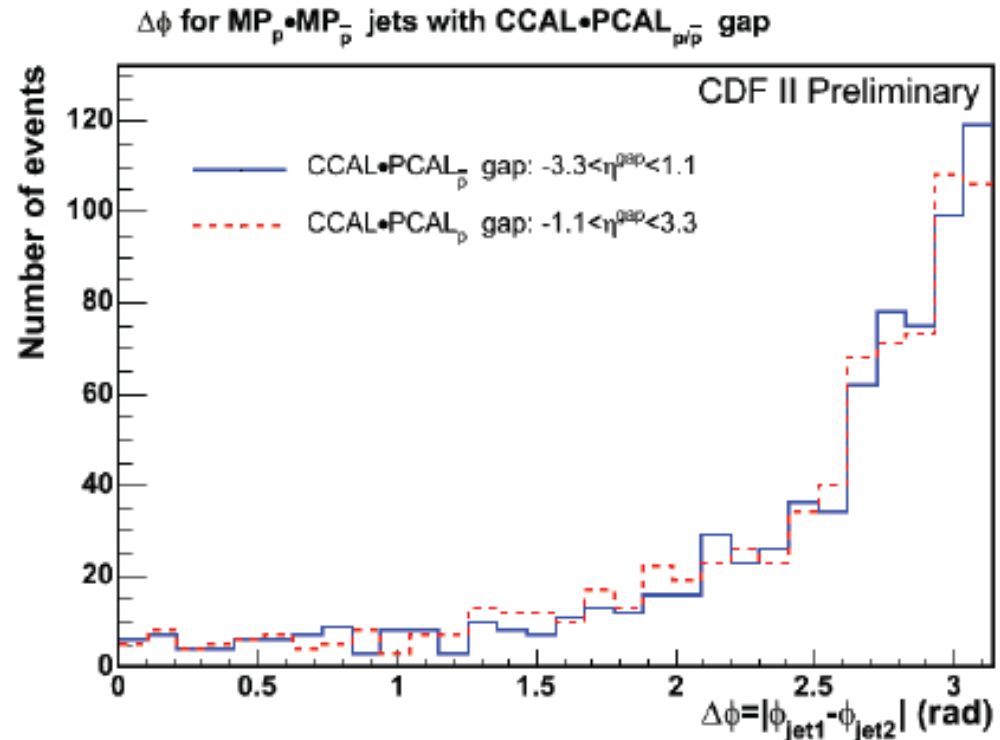
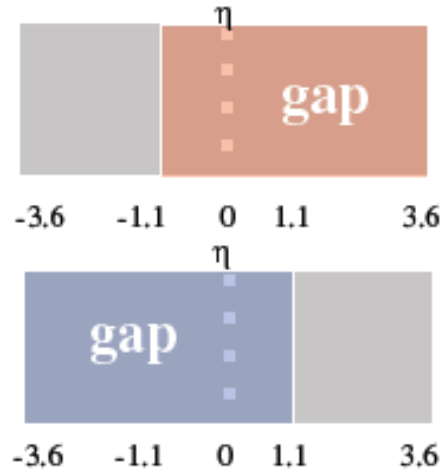
The number of CCAL towers with  $E_T > 200 \text{ MeV}$  is suppressed by the TOF veto

# CCAL gap



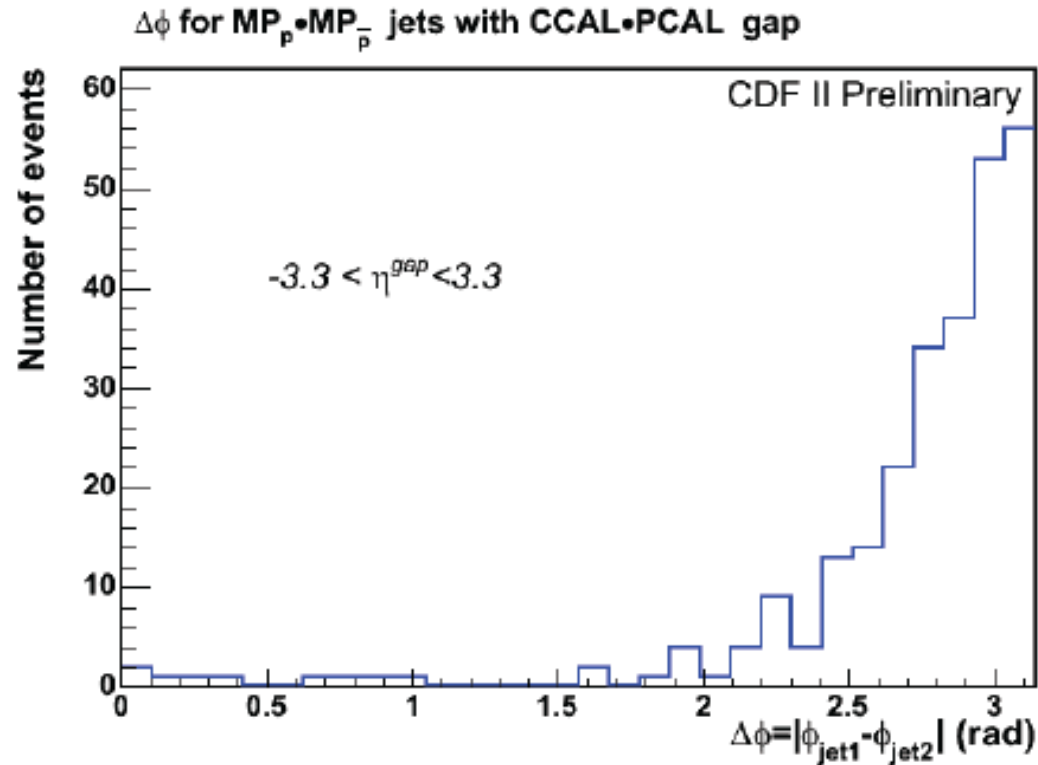
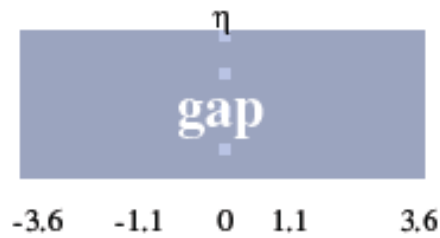
*The distribution of  $\Delta\phi$  for  $MP_p \bullet MP_{pbar}$  jets of  $E_T > 2$  GeV with a gap in the central calorimeter (CCAL). The events at low values of  $\Delta\phi$  are presumed to be due to an imbalance caused by the  $E_T$  of  $PCAL_p$  and  $PCAL_{pbar}$  towers from the underlying event.*

# CCAL • PCAL<sub>p</sub>/<sub>pbar</sub> Gap



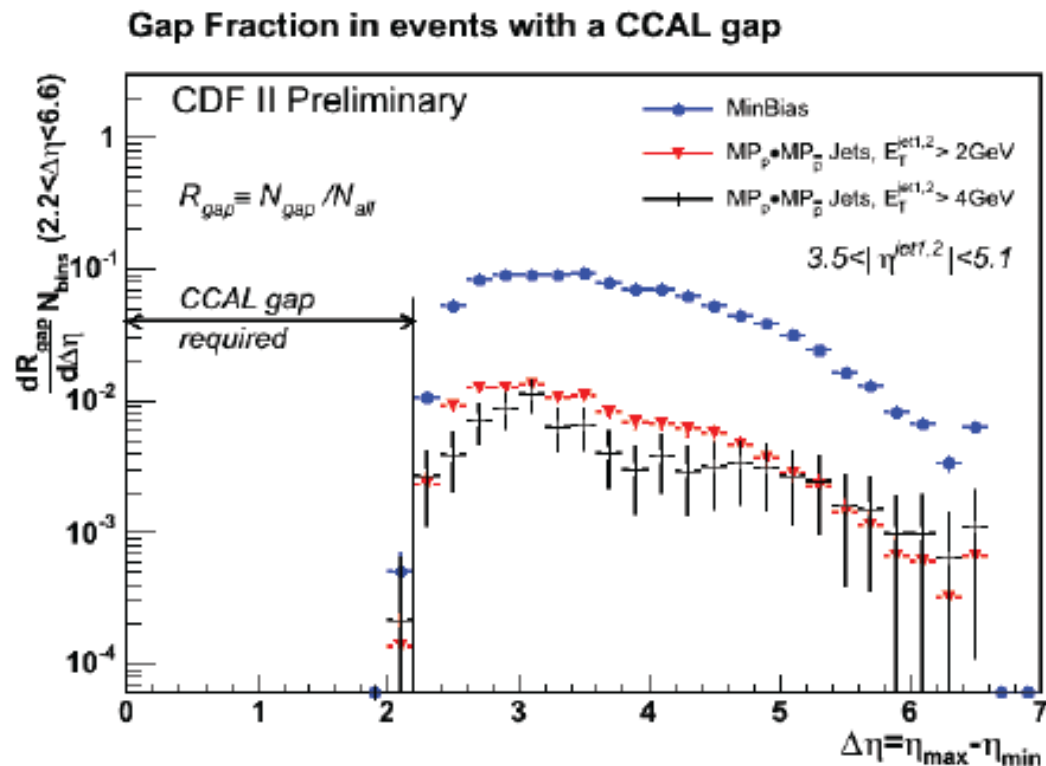
*The distribution of  $\Delta\phi$  for  $MP_p \bullet MP_{pbar}$  jets of  $E_T > 2$  GeV with a gap in the central calorimeter (CCAL) plus a gap in one of the Plug calorimeters PCAL(p) or PCAL(pbar). The events at low values of  $\Delta\phi$  are presumed to be due to the imbalance caused by the  $E_T$  of PCAL(p) or PCAL(pbar) towers from the underlying event. The agreement between the two distributions indicates that detector and beam conditions effects are similar between positive and negative eta-values.*

# CCAL • PCAL<sub>p</sub> • PCAL<sub>pbar</sub> Gap



*The distribution of  $\Delta\phi$  for  $MP_p \bullet MP_{pbar}$  jets of  $E_T > 2$  GeV with a gap in the central calorimeter (CCAL) plus a gap in both Plug calorimeters - PCAL(p) and PCAL(pbar).*

# Gap fraction in CCAL gap Events



The distribution of the gap fraction  $R_{\text{gap}} = N_{\text{gap}} / N_{\text{all}}$  vs  $\Delta\eta$  for MinBias ( $CLC_p \bullet CLC_{pbar}$ ) and MiniPlug jet events ( $MP_p \bullet MP_{pbar}$ ) of  $E_{T(\text{jet}1,2)} > 2 \text{ GeV}$  and  $E_{T(\text{jet}1,2)} > 4 \text{ GeV}$ .

**The distributions are similar in shape within the uncertainties.**

# SUMMARY

## □ Introduction

- Renormalization,  $M^2$ -scaling, survival probability
- Diffractive and non-diff. proton PDFs similar

## □ Diffractive dijets – $x_{Bj}$ , $Q^2$ and $t$ -dependence

## □ Diffractive W/Z with RPS data

- W diffractive fraction in agreement with Run I
- W and Z diffractive fractions are equal within error

## □ Exclusive dijets and exclusive Higgs @ LHC

## □ Exclusive di-lepton and di-photon measurements

## □ Exclusive Z production limit

## □ Central rapidity gaps

- Gap fraction dependence on width and  $\eta$ -position of gap for hard / soft triggers at  $|\eta| > 4$ 
  - ➔ distributions shapes similar for hard / soft triggers
  - ➔ hard-scale fractions suppressed by factor of or  $\sim 10$ .



thank you



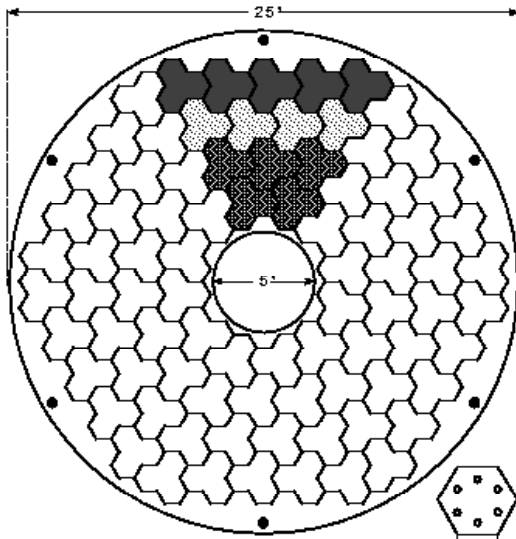
# BACKUP

Measurements w/ the MiniPlugs

Dynamic Alignment of RPS Detectors

$E_{\text{jet}}^T$  Calibration

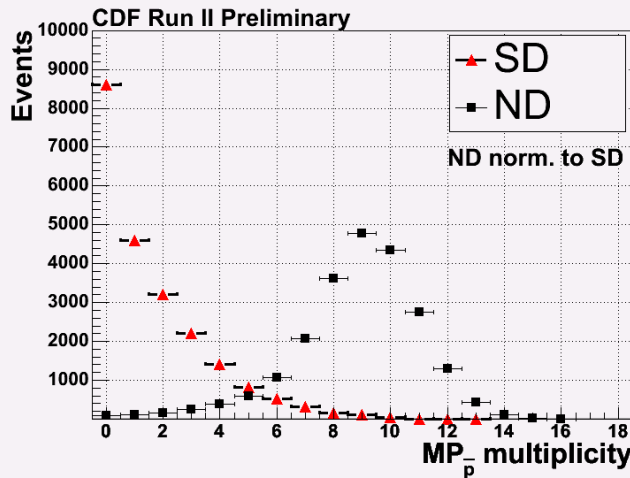
# Measurements w/ the MiniPlugs



← MP TOWER  
STRUCTURE

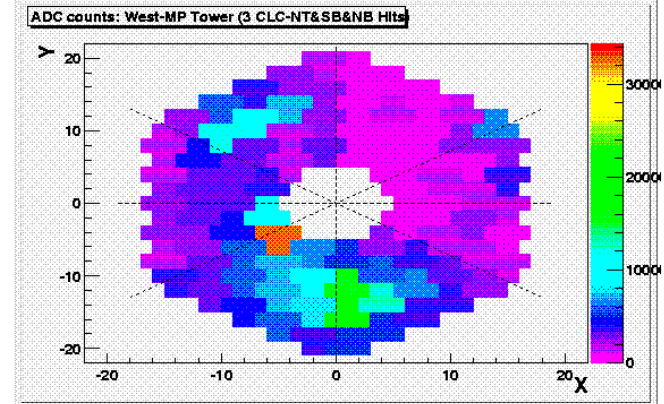
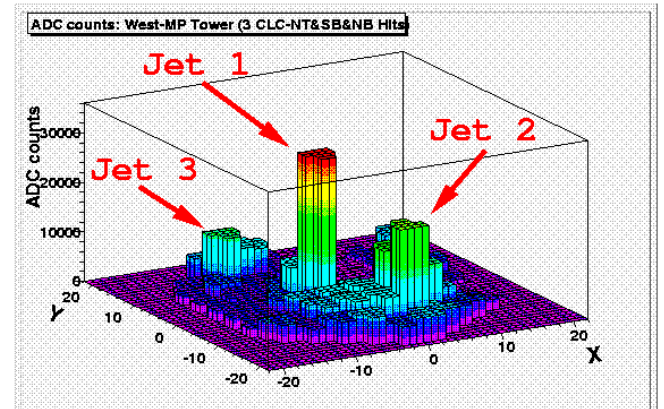
→ MULTIPLICITY  
@ POSITION

→ ENERGY



$$\xi_{CAL} = \frac{\sum_i E_T^i e^{-\eta_i}}{\sqrt{s}}$$

NIM A 430 (1999)  
NIM A 496 (2003)  
NIM A 518 (2004)



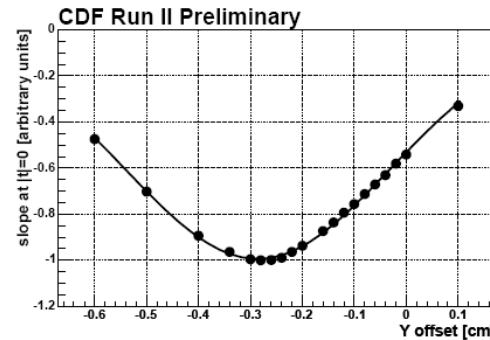
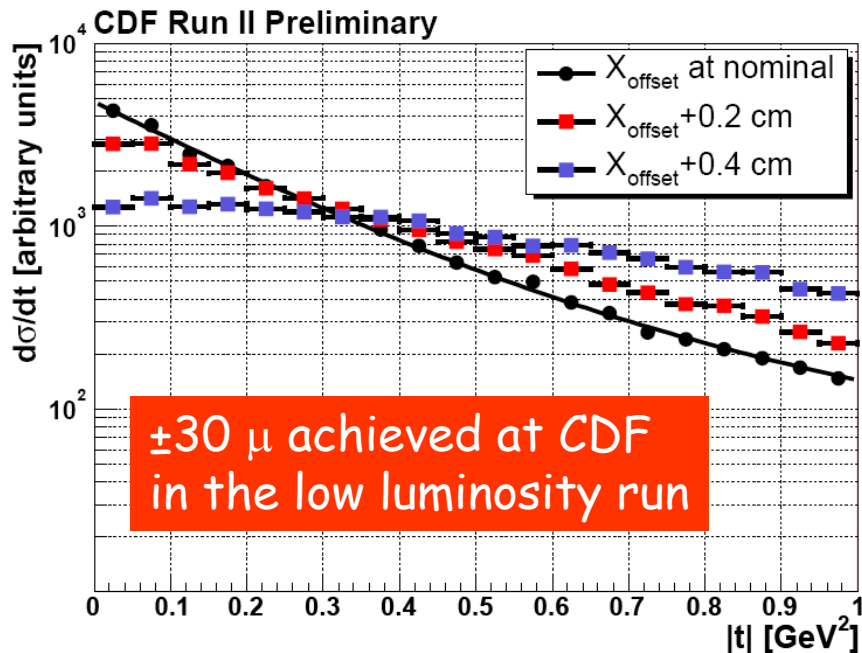
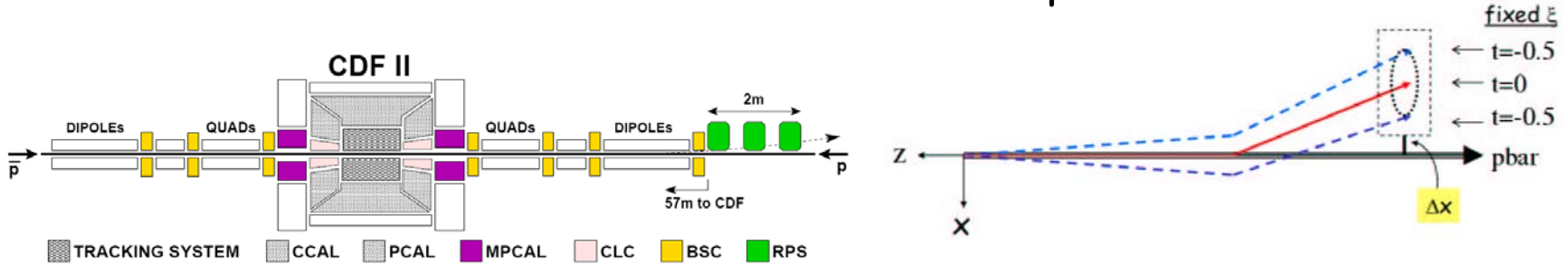
ADC counts in MiniPlug towers in a  $p\bar{p}$  event at 1960 GeV.

- “jet” indicates an energy cluster and may be just a hadron.
- 1000 counts ~ 1 GeV

Multiplicity of SD and ND events

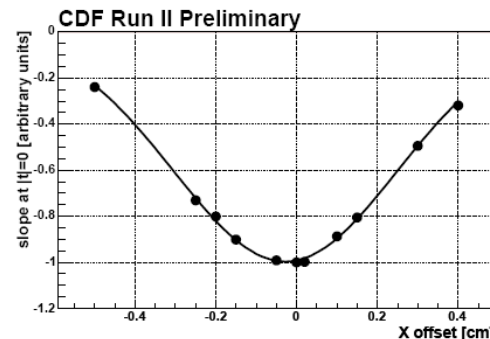
# Dynamic Alignment of RPS Detectors

Method: iteratively adjust the RPS X and Y offsets from the nominal beam axis until a maximum in the b-slope is obtained @  $t=0$ .



Limiting factors

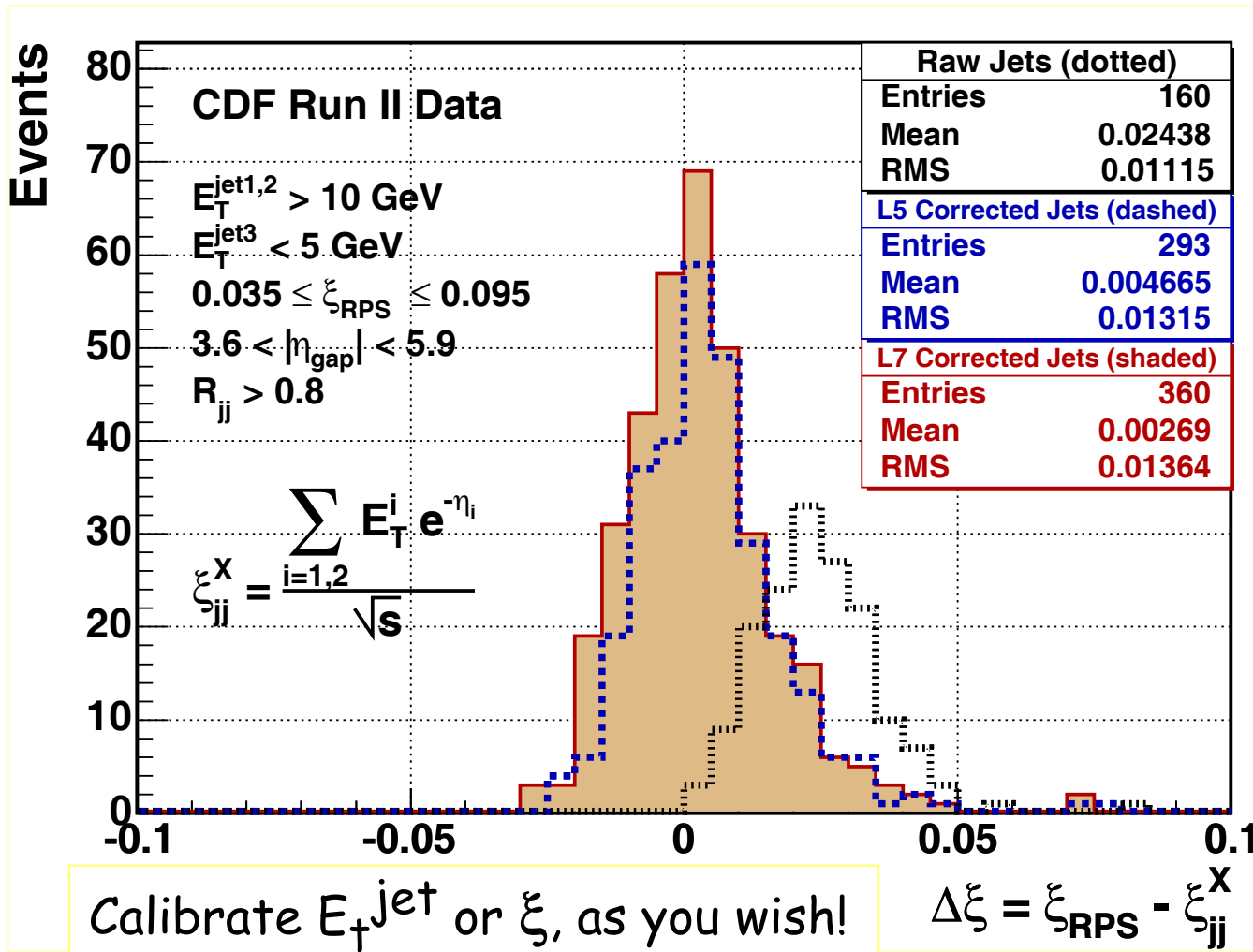
- 1-statistics
- 2-beam size
- 3-beam jitter



@ CDF  
w/lowlum data  
 $\pm 30 \mu\text{m}$

# $E_{T}^{\text{jet}}$ Calibration

→ use RPS information to check jet energy corrections ←



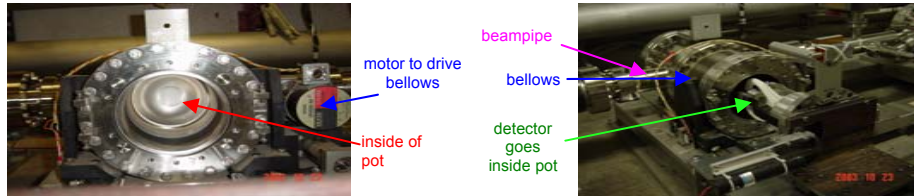
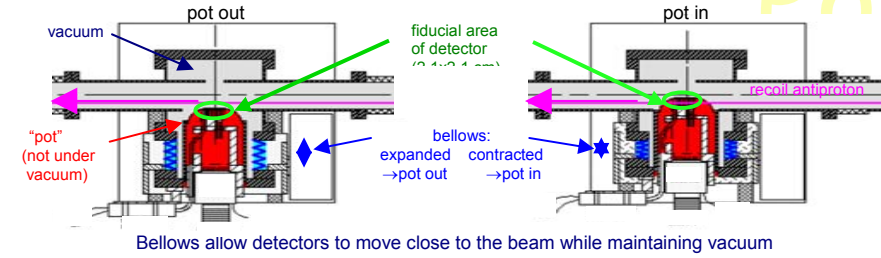




# The Roman-Pot Detectors at CDF



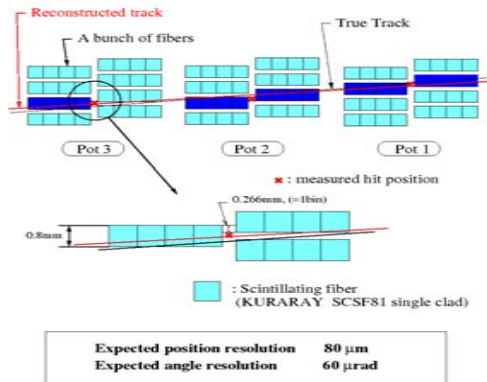
## Concept of a Roman Pot



## Roman-Pot Detector Design – by The Rockefeller University

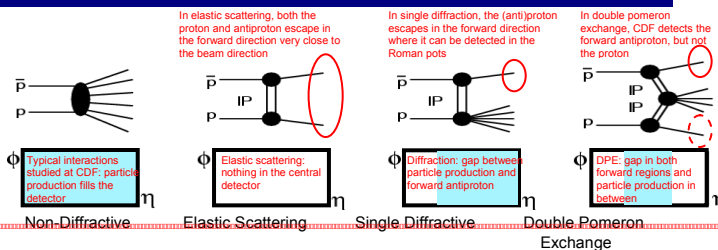
The three Roman pots each contain detectors consisting of:

- Trigger scintillating counter 2.1x2.1x0.8 cm<sup>3</sup>
- 40 X + 40 Y fiber readout channels
  - Each consists of 4 (→ bigger signal) clad scintillating fibers 0.8x0.8 mm<sup>2</sup> (new technology at the time)
  - X, Y each have 2 rows of 20 fibers spaced 1/3 fiber width apart for improved position resolution (three times better than with a single row)



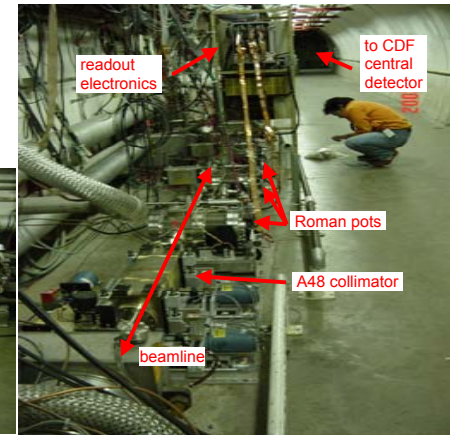
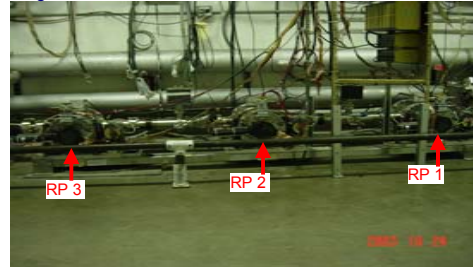
## Physics Using the Roman-Pot Detectors

- The Roman-pot detectors are used to study diffractive interactions
- Elastic scattering was measured by CDF in 1988-1989 using Roman pots (not those described here) in both the proton and antiproton direction

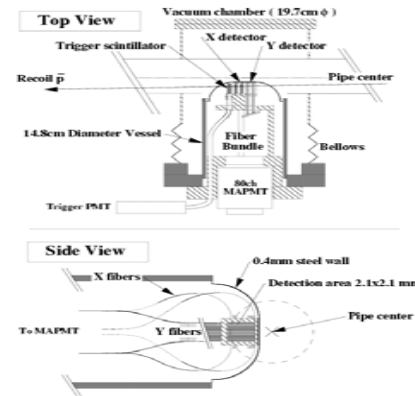


## In the Tevatron Tunnel

CDF had three Roman pots (RP1, RP2, RP3) located 57m downstream of the interaction point along the antiproton beam direction. They were used to detect antiprotons which underwent a "diffractive" interaction and were scattered in a direction very close to that of the original beam.



## CDF "Tokyo"-Pot Detectors – Built by the University of Tsukuba, Japan



## Path of the Antiproton through the Tevatron Magnets

- Dipole magnets bend recoil antiprotons which have lost momentum towards the inside of the Tevatron ring, into the Roman pots
- Knowledge of the beam optics, the collision vertex position, and the antiproton track position and angle in the Roman-pot detectors are used to reconstruct the kinematics of the diffractive antiproton

